



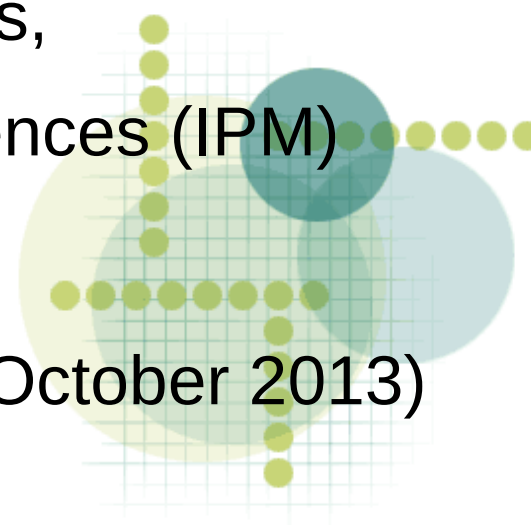
# Hadronic Search for SUSY with $MT_2$ variable

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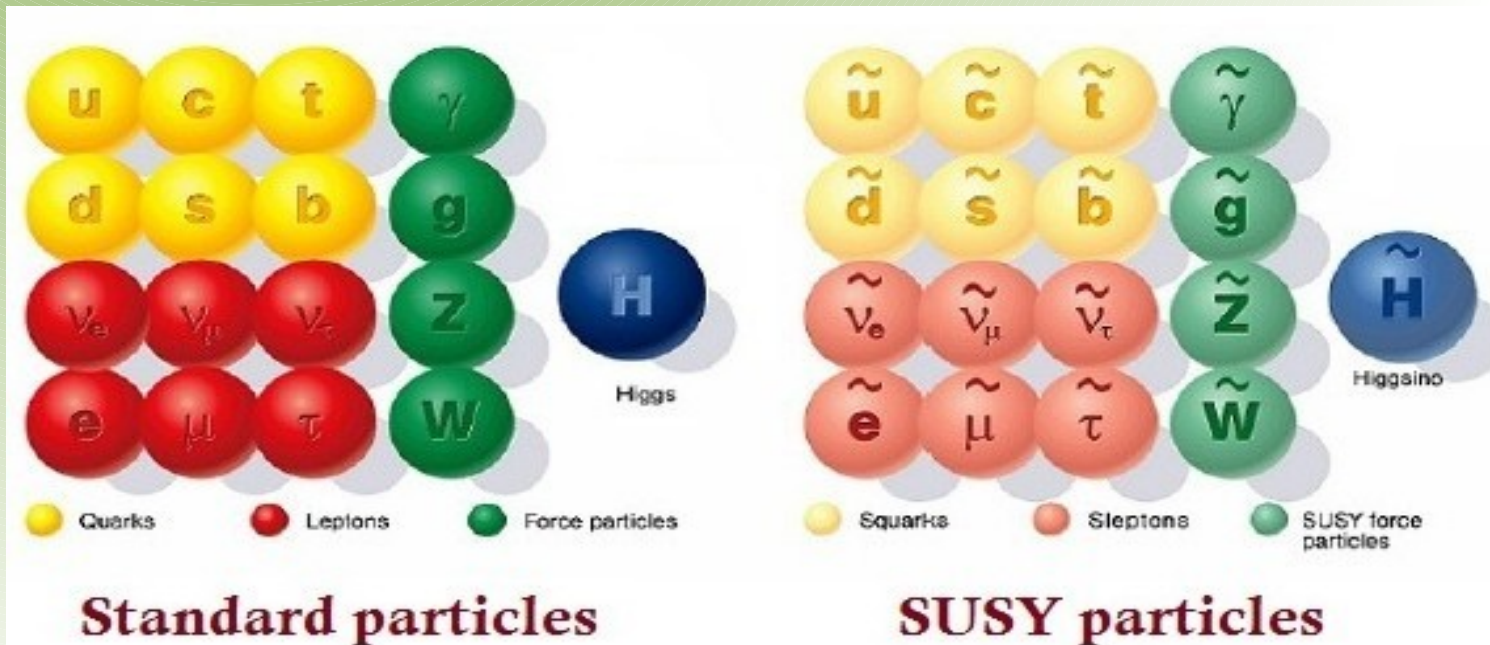
# Outline

- Supersymmetry and Its Motivations
- Why Hadronic search? Why  $MT_2$ ?
- From  $MT$  to  $MT_2$
- Interpretation of  $MT_2$
- Search Strategy
- Backgrounds Estimation
- Results
- Conclusion



# Supersymmetry and Its Motivations

- SuperSYmmetry (SUSY) is a symmetry between fermions  $\leftrightarrow$  bosons,  
 $Q |\text{boson}\rangle = |\text{fermion}\rangle$ ,  $Q |\text{fermion}\rangle = |\text{boson}\rangle$



- SUSY: double number of particles (MSSM)
- Spin of SUSY particles differs by 1/2 a unit
- Must be broken, 105 free parameters due to SUSY breaking
- mSUGRA/cMSSM has 5 free parameters:  $A_0$ ,  $\tan\beta$ ,  $\text{sgn}(\mu)$ ,  $m_0$ ,  $m_{1/2}$ .

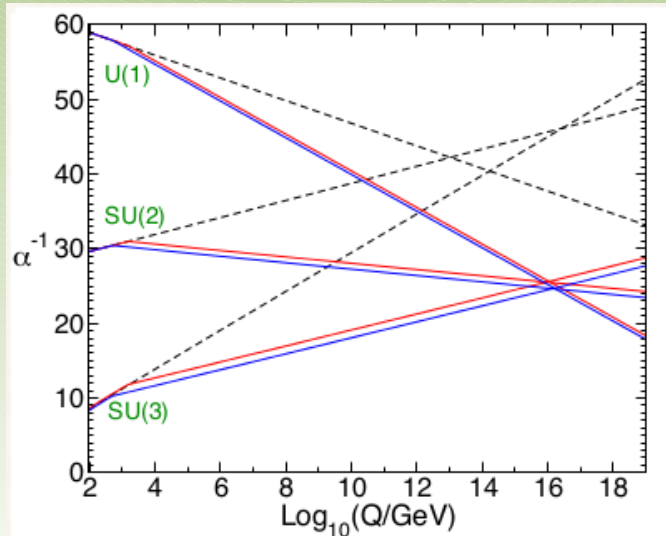


# Supersymmetry and Its Motivations

- SUSY is very interesting, as it:

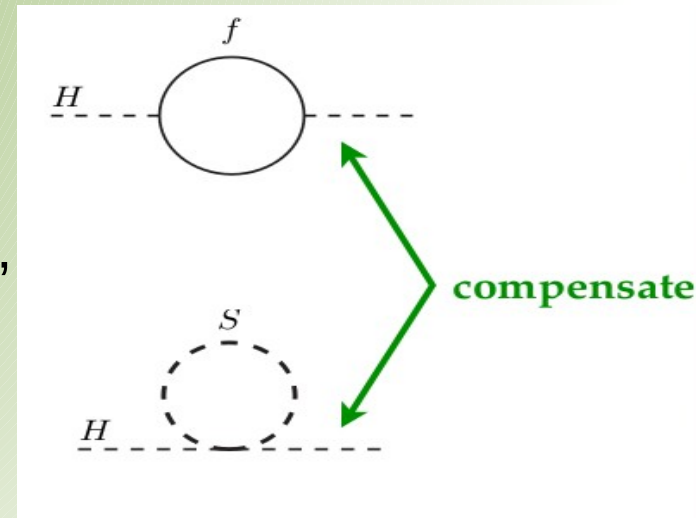
1. gives a

“Solution to the Hierarchy Problem”



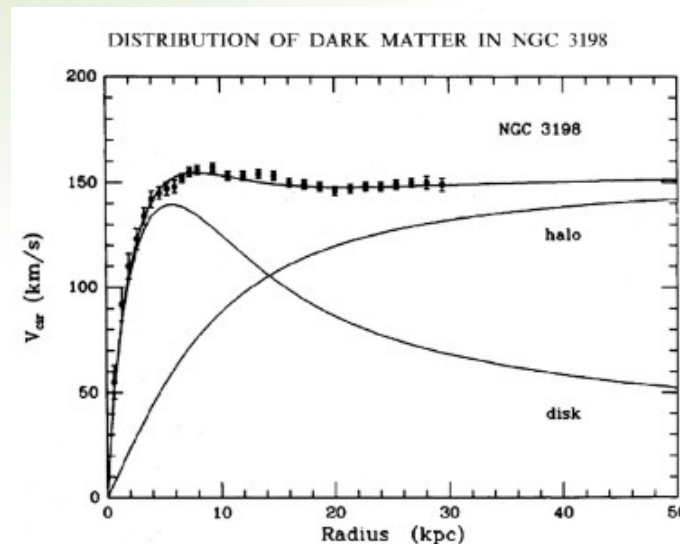
3. provides a

“Dark Matter candidate”



2. facilitates a

“Gauge Couplings Unification”



Astrophysical Journal,  
vol. 295, Aug. 15, 1985,  
p. 305-313.


# Why Hadronic search? Why MT2?

If SUSY ,conserving R-parity, comes with a Dark Matter candidate:


- SUSY particles must be produced in pairs and the lightest sparticle (LSP) as a Dark matter candidate is stable.
- Colored SUSY particles cascade down to LSP with emission of jets and sometimes leptons, therefore in search for SUSY, MET and HT are very useful variables:

$$H_T = \sum_{jets} |\vec{p}_T|$$

$$E_T^{miss} = \left| - \sum_{particles} \vec{p}_T \right|$$

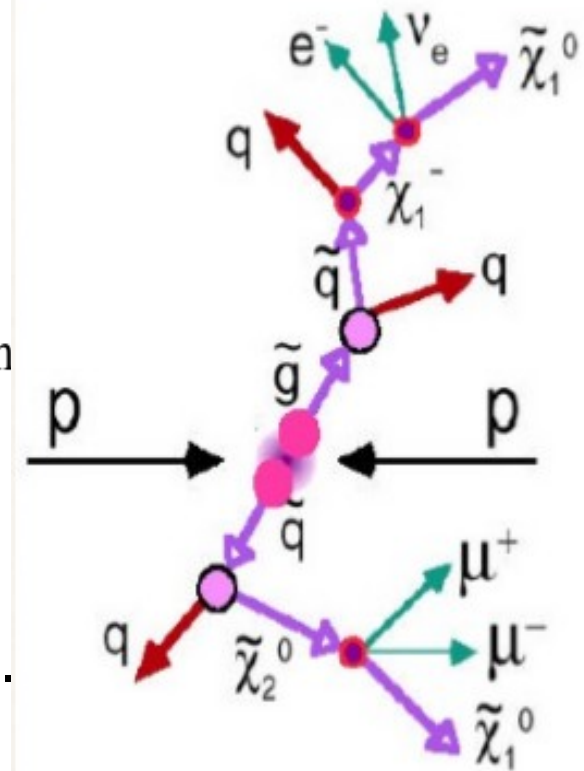


Hadronic activity



Missing transverse momentum

- Classical SUSY searches are based on these two variables but it can be defined a clever kinematical variable reflecting these properties which is “MT2”.



# From MT to MT2

In  $W \rightarrow e\nu$  decay, **transverse mass MT** has an endpoint at the true W-mass:

$$m_W^2 = m_l^2 + m_\nu^2 + 2(E_T^l E_T^\nu \cosh \Delta\eta - \mathbf{p}_T^l \cdot \mathbf{p}_T^\nu) \geq$$

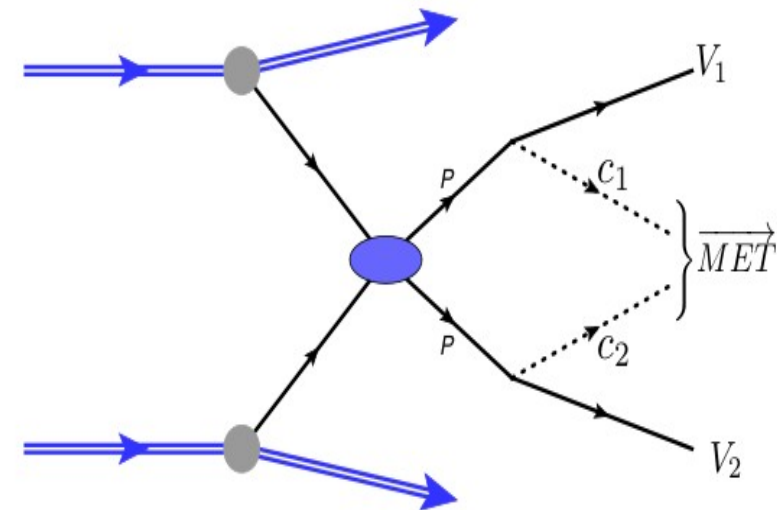
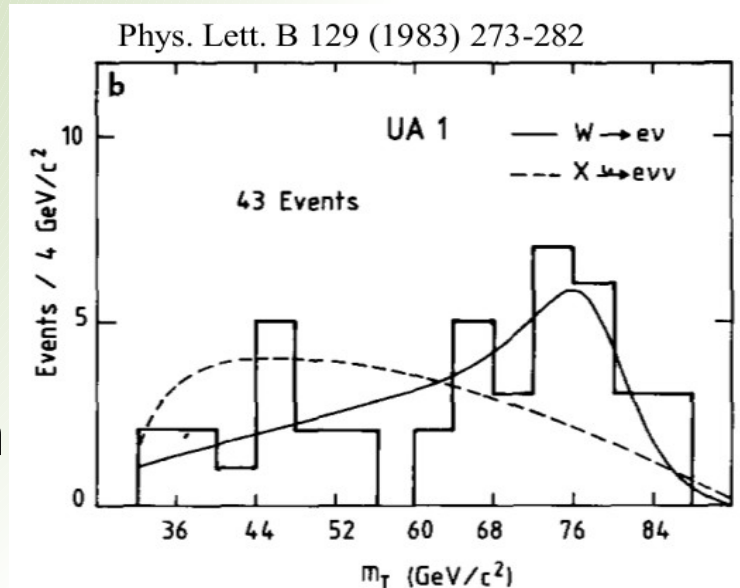
$$m_T^2 = m_l^2 + m_\nu^2 + 2(E_T^l E_T^\nu - \mathbf{p}_T^l \cdot \mathbf{p}_T^\nu)$$

At the LHC, assuming R-parity conservation, **SUSY** events give rise to two decay chains with an unobserved child (c1 and c2) at each end.

The “**stransverse**” mass **MT2**: extension of MT for the **SUSY** case of two unobserved particles:

$$M_{T2}(m_c) = \min_{p_T^{c(1)} + p_T^{c(2)} = p_T^{miss}} \left[ \max \left( m_T^{(1)}, m_T^{(2)} \right) \right]$$

If  $M_c$  were known, the endpoint of MT2 would correspond to the parent mass  $M_p$ .





# Interpretation of MT2

- In case of **no initial state radiation (ISR)** and **zero masses**:

$$M_{T2}^2 = 2 p_T^{(1)} p_T^{(2)} (1 + \cos \phi_{1,2})$$

which  $p_T(i)$  is the transverse momenta of the visible systems.

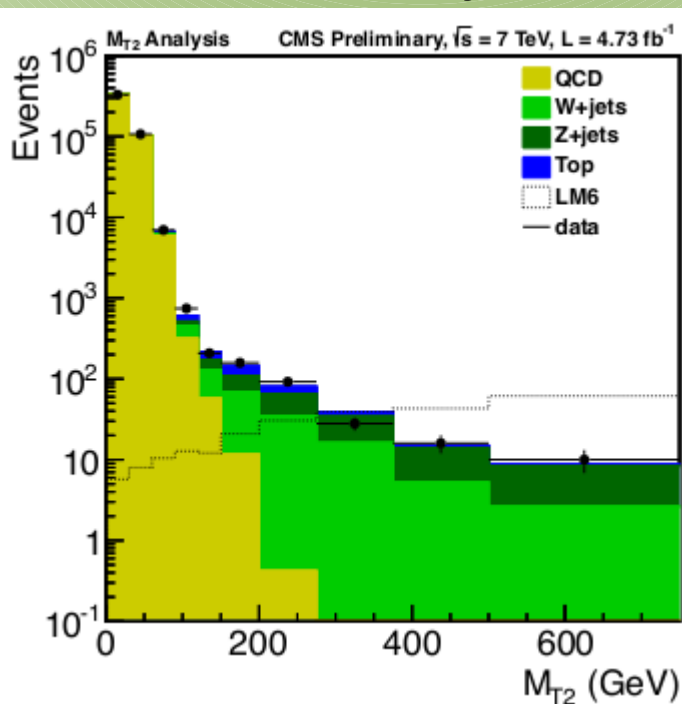
- MT2 = MET** for symmetric systems,  $p_T(1)=p_T(2)$

$$(E_T^{\text{miss}})^2 = (p_T^{(1)} - p_T^{(2)})^2 + 2 p_T^{(1)} p_T^{(2)} (1 + \cos \phi_{12})$$

- MT2 = 0** GeV for **back-to-back systems**, like QCD di-jets events.
- MT2 is similar to MET in signal region, but **more robust against jet energy mismeasurements than MET**. Therefore QCD multijet events accumulated at low MT2.  
Events with real MET can have large values in MT2.
- Multijet events are divided into a 2 pseudo-jets topology using a **hemisphere algorithm**.

# Search Strategy

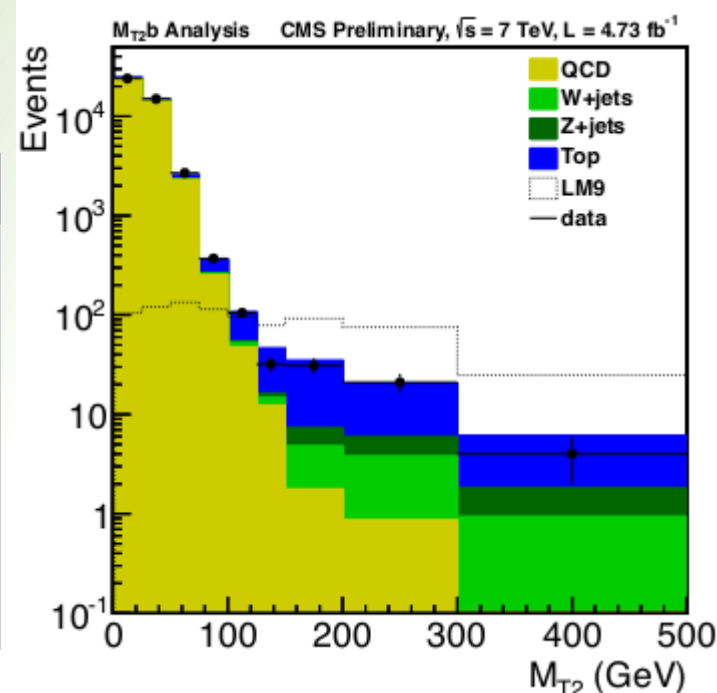
- Search in fully hadronic final states with 2011 pp collision data at 7 TeV collected by the CMS detector, corresponding to 4.73 fb<sup>-1</sup>.



## Baseline selection

$$HT \geq 750 \text{ GeV}$$

- electron and muon veto
- filters against detector noise



## MT2 analysis

- At least 3 jets
- $\min\Delta\phi(\text{jets}, \text{MET}) > 0.3$
- $MT2 > 150 \text{ GeV}$

sensitive to high/medium  
squark and gluino masses

sensitive to light gluinos  
with heavy squarks

## MT2b analysis

- At least 4 jets
- At least one jet b-tagged
- $\min\Delta\phi(\text{j1-j4}, \text{MET}) > 0.3$
- $MT2 > 125 \text{ GeV}$



# Backgrounds Estimation

- **QCD multijets** have no genuine MET → **small MT2**. MC predicts that the signal regions are QCD free.
- Mismeasured jets can lead to larger MT2. These jets are aligned with MET. Use **correlation of  $\min\Delta\phi(\text{jets}, \text{MET})$  and MT2** to conservatively estimate the QCD contamination in the signal region from data.

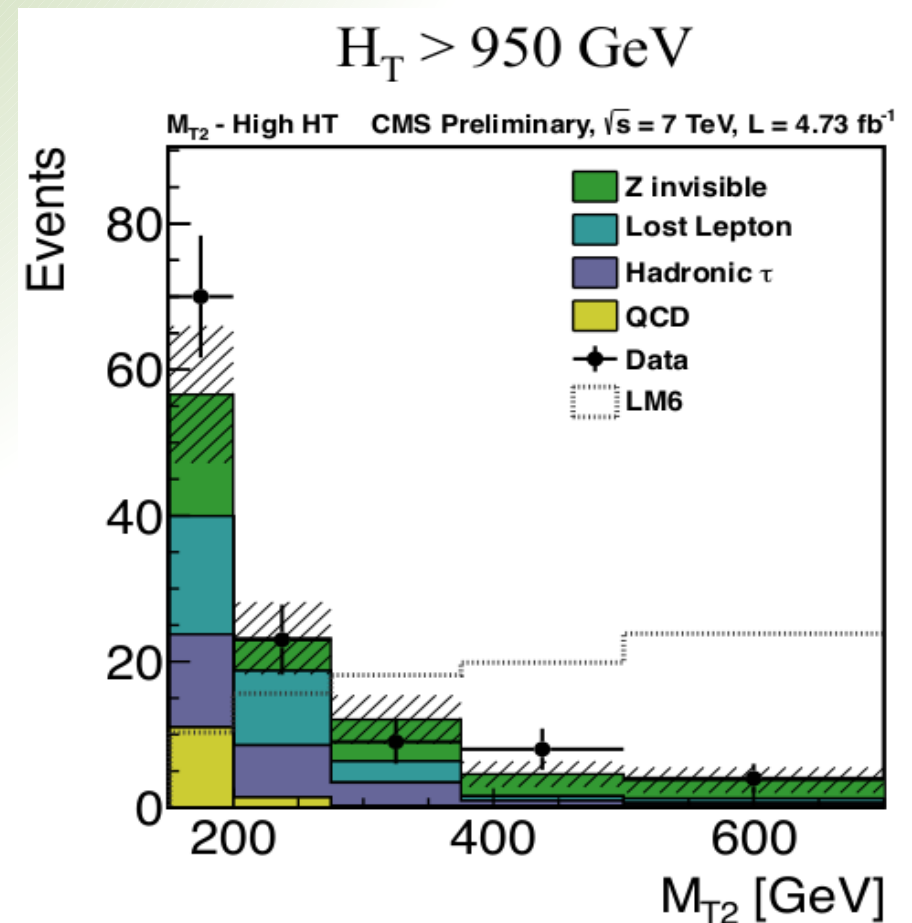
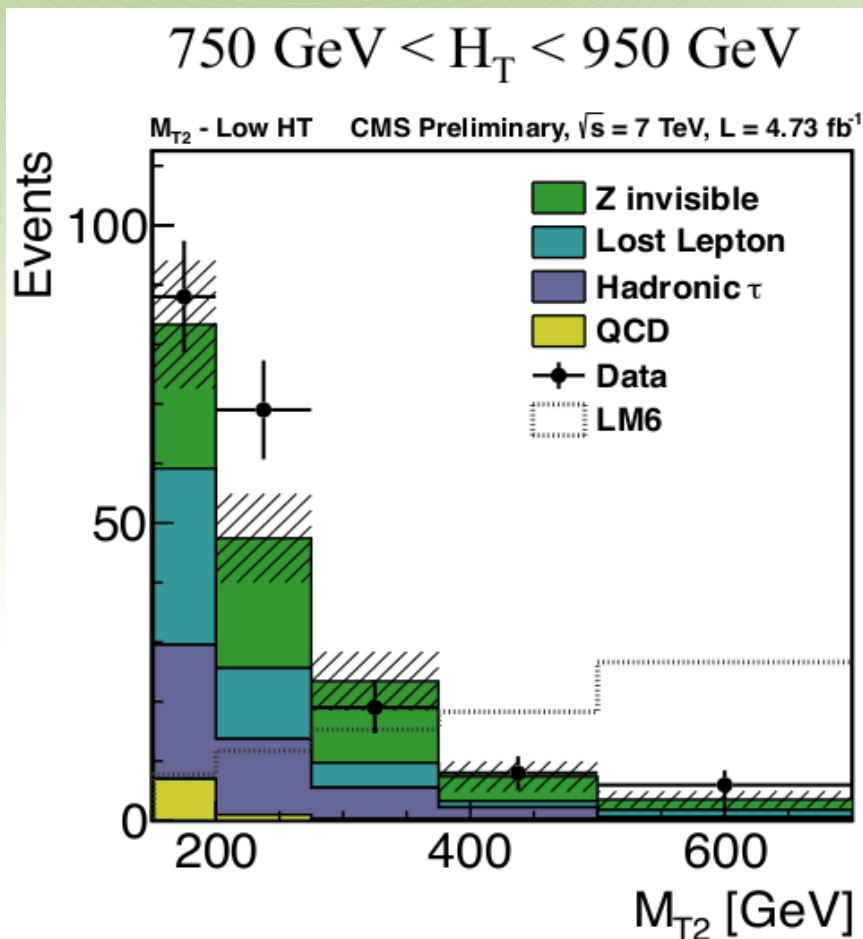
- **Z bosons decaying into neutrinos** are signal like with large MET.
- This background is predicted by using  **$W(\rightarrow \mu\nu)$ +jets and photon + jets** events.
- In both cases the visible vector boson pT is added to the MET to mimic  $Z(\rightarrow \nu\nu)$ +jets events.

- **Leptonic W+jets and Top+jets** events have real MET. Largely reduced due to lepton veto.
- Enter to signal region if the charged lepton is not reconstructed or out of acceptance (= is lost).
- Remaining (**lost lepton**) background estimated in EWK control region from the number of events with a found lepton in data and corrected for the probability to lose a lepton.

**Robust estimation of the SM backgrounds contribution to all signal bins.**

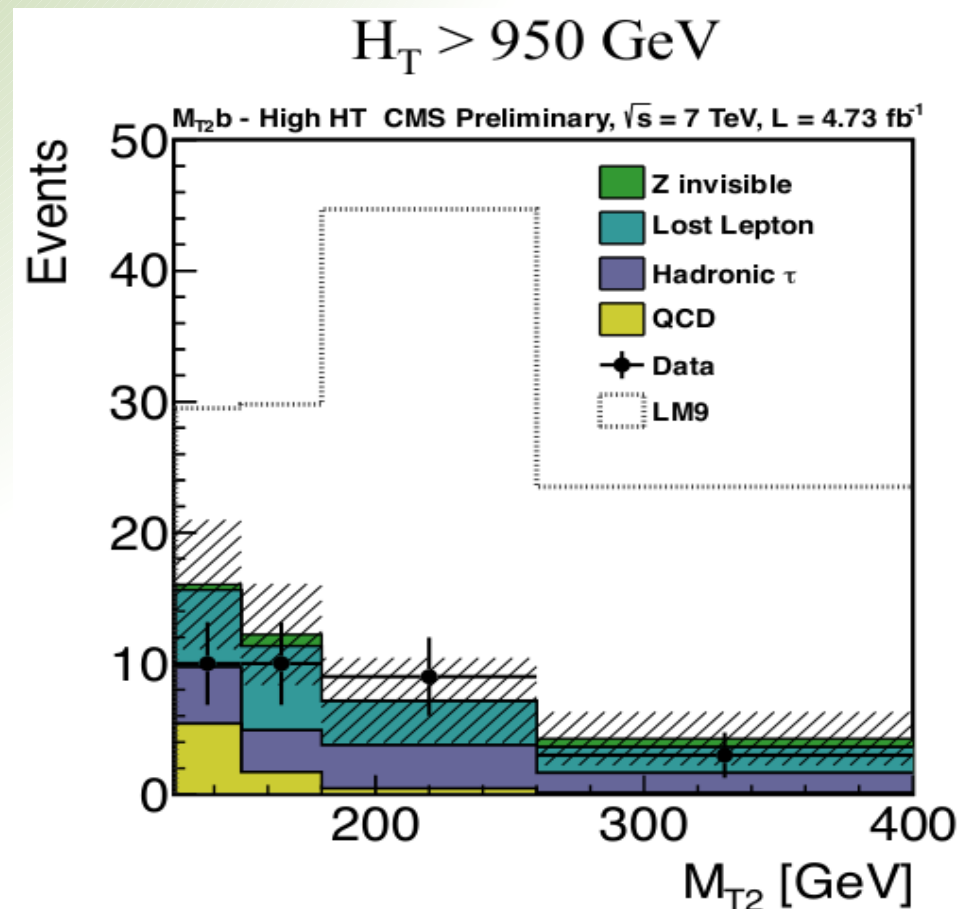
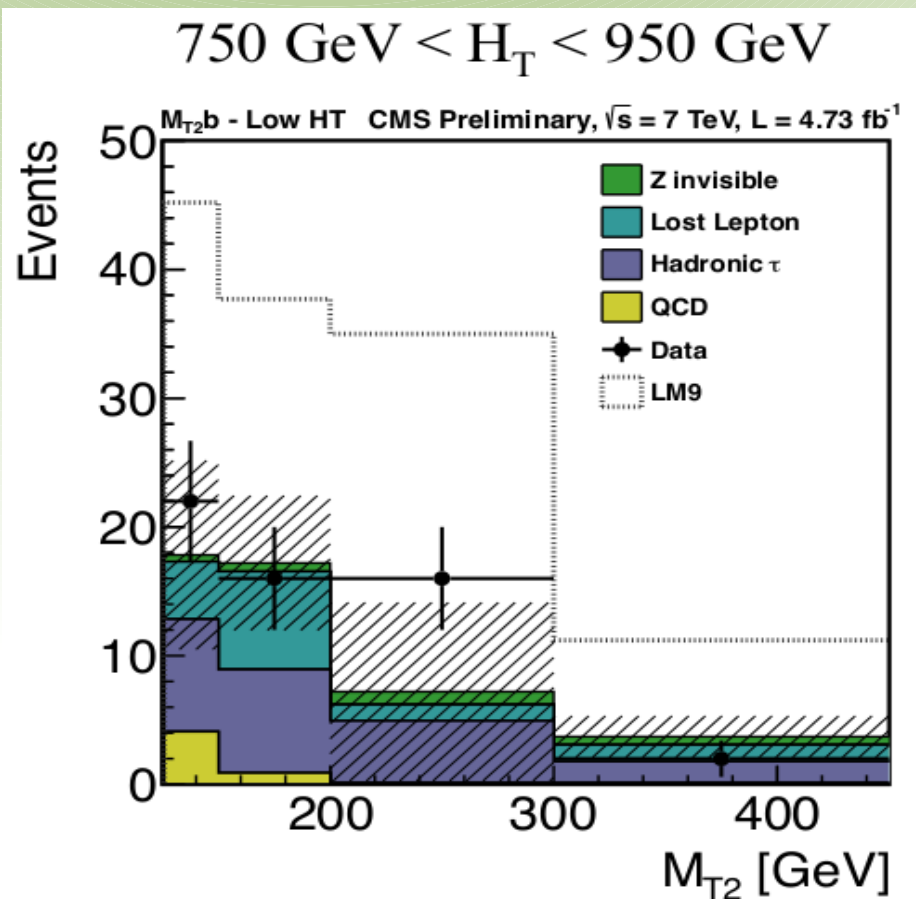
# Backgrounds Estimation

- The backgrounds estimation for the MT2 analysis are summarized here.
- Shaded region is uncertainty on the background estimation.
- The observed data are shown.
- A possible SUSY signal is overlaid to show sensitivity of search region.



# Backgrounds Estimation

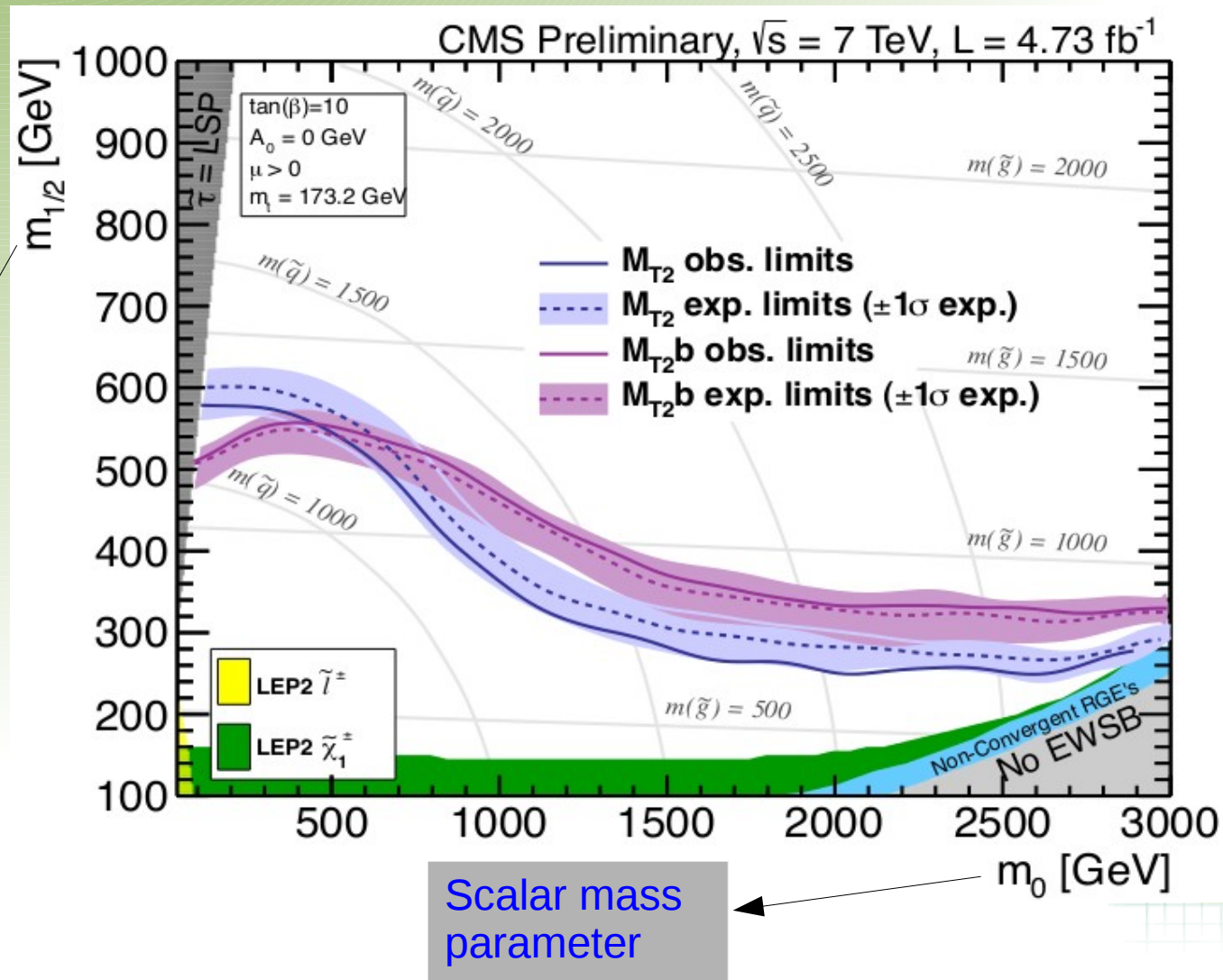
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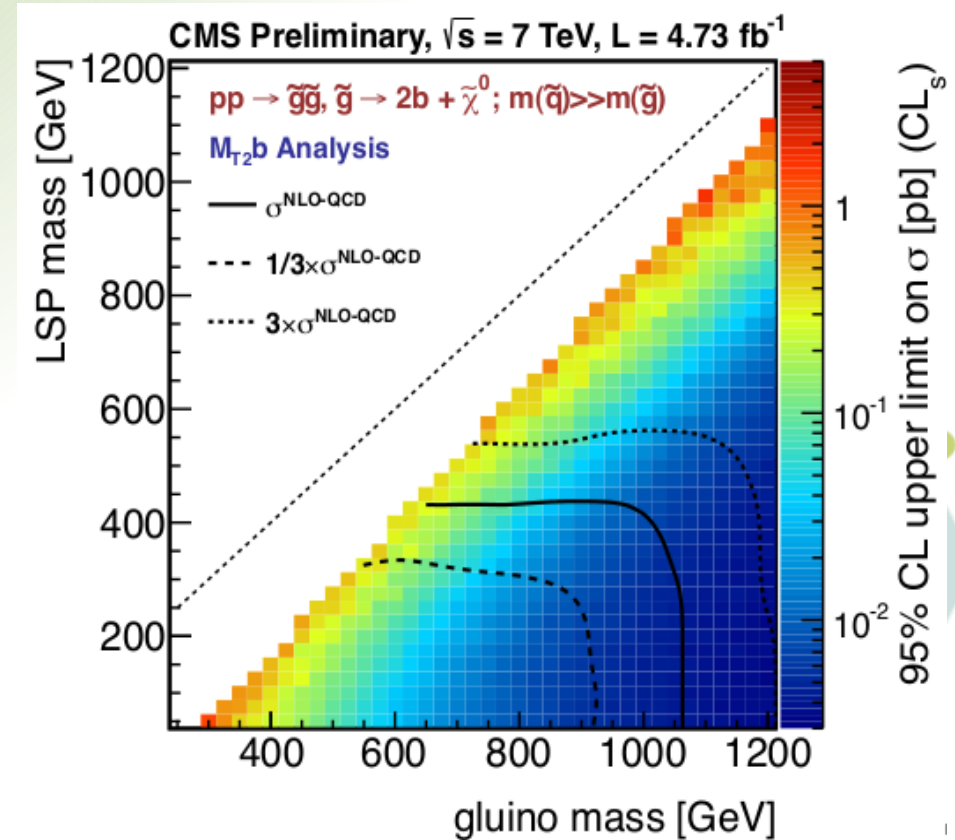
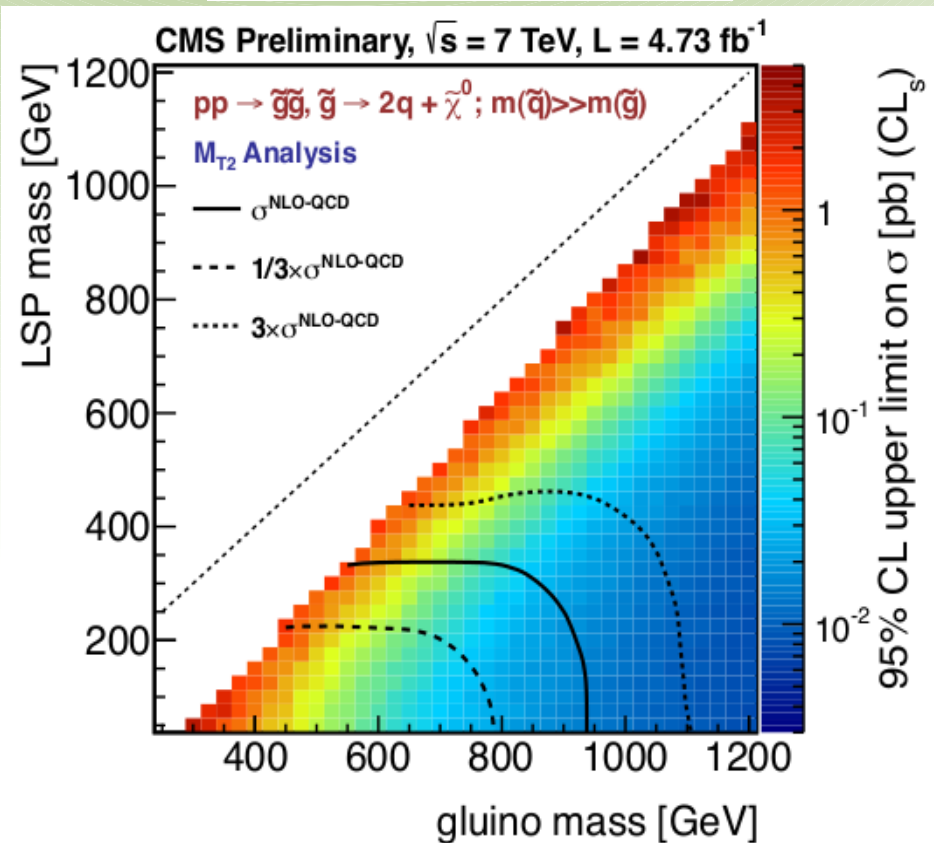
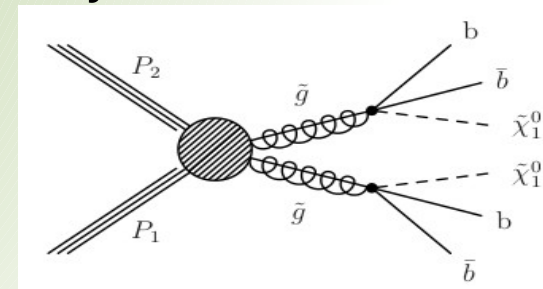
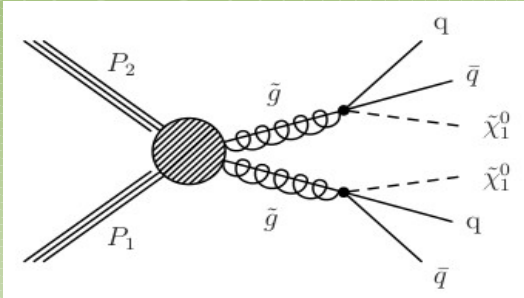
# Results

- The results are interpreted in a full SUSY model constrained to five parameters (mSUGRA/cMSSM).
- In the plane below three of those parameters are fixed:  $A_0 = 0$ ,  $\tan\beta = 10$ ,  $\mu > 0$



# Results

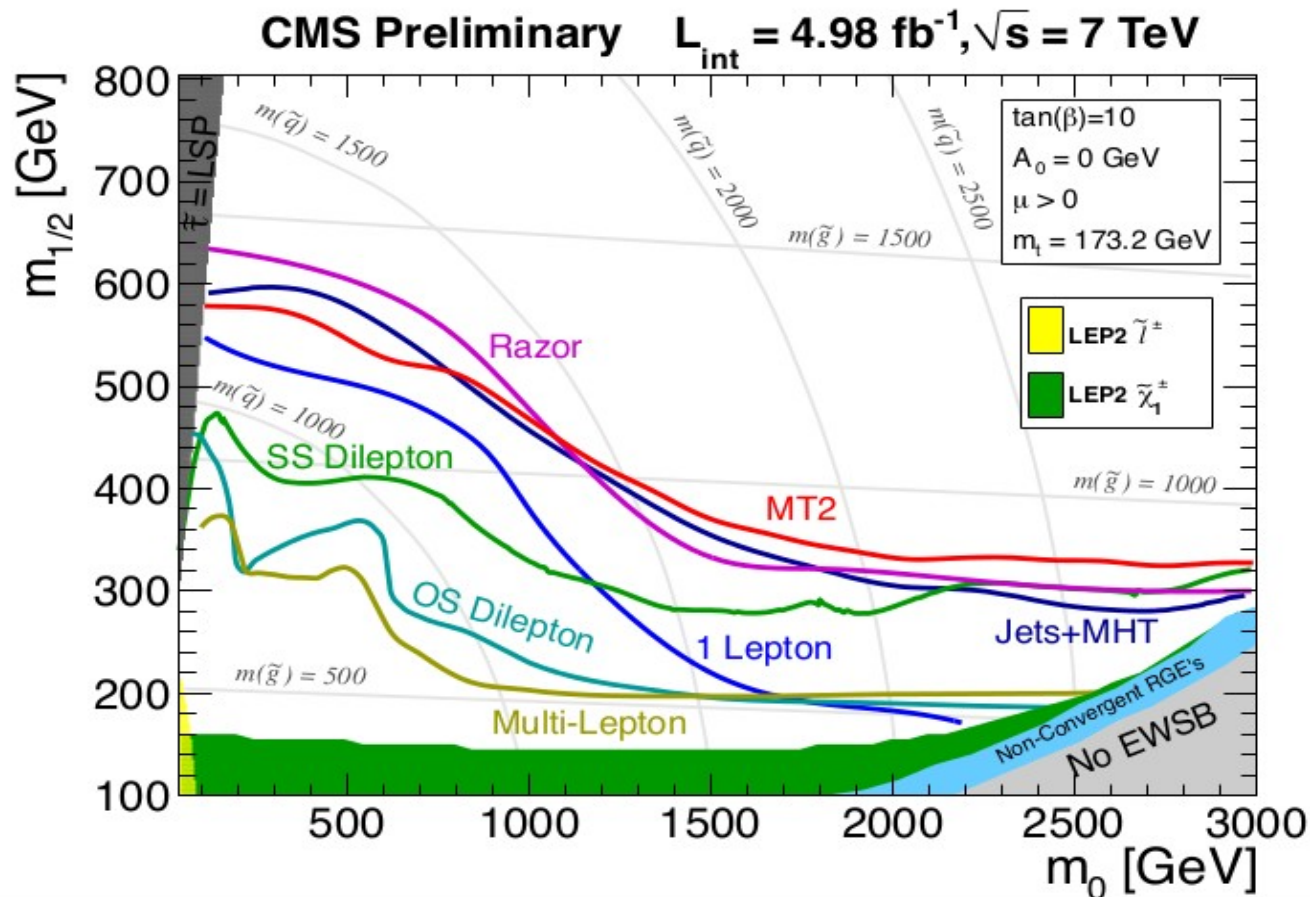
- The analyses are also interpreted in simplified models.
- Models are reduced to one SUSY decay chain only.





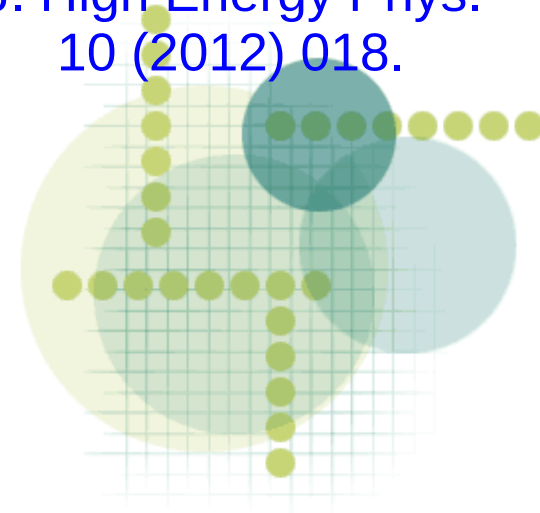
# Conclusion

- A search for supersymmetry in fully hadronic final states with 2011 pp collision data collected by the CMS detector has been performed.
- No excess over the SM predicted background has been found.
- Limits in various signal model spaces have been set. Absolute mass limits in the mSUGRA/cMSSM scenario for  $A_0 = 0$ ,  $\tan\beta = 10$ ,  $\mu > 0$  are found to be  $m(\text{squark}) > 1160 \text{ GeV}$  and  $m(\text{gluino}) > 860 \text{ GeV}$ , and  $m(\text{squark}) = m(\text{gluino}) > 1200 \text{ GeV}$  assuming equal squark and gluino masses.



For more details:

CMS-PAS SUS-12-002;  
arXiv:1207.1798;  
Published in :  
J. High Energy Phys.  
10 (2012) 018.





Thank you.



*"One day, all of these will be supersymmetric phenomenology papers."*