### Supersymmetric continuous spin gauge theory

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Tehran

#### Based on:

M. Najafizadeh, "Supersymmetric Continuous Spin Gauge Theory"

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#### Outline:

- Supersymmetry (SUSY)
- Continuous Spin Particle (CSP)
- Continuous spin gauge field theory
- SUSY CSP

convention:  $p^2 = 0$ ,  $\mathcal{N} = 1$ , d = 4



- a symmetry between fermions and bosons
- SUSY particles have not yet been observed
- zero-point energy
- supercharge Q:  $Q \mid \text{boson} > = \mid \text{fermion} > Q \mid \text{fermion} > = \mid \text{boson} >$
- supermultiplet: (boson, fermion)where # bosonic d.o.f = # fermionic d.o.f
  - $(0, \frac{1}{2}) \longrightarrow \text{chiral multiplet (Wess-Zumino, 1974)}$
  - $(1, \frac{1}{2}) \longrightarrow \text{gauge multiplet}$
  - $(2, \frac{3}{2}) \longrightarrow \text{gravity multiplet}$
  - $(s, s+\frac{1}{2}) \longrightarrow \text{half-integer spin multiplet (Curtright, 1979)}$  $(s+\frac{1}{2}, s+1) \longrightarrow \text{integer spin multiplet}$  (Curtright, 1979)

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To supersymmetrize a theory, given by a sum of bosonic and fermionic actions

$$S_{\scriptscriptstyle SUSY} = S_b \left[ \phi \right] + S_f \left[ \psi \right]$$

one should find "SUSY transformations":  $\begin{cases} \delta \phi = \overline{\epsilon} (\cdots) \psi \\ \delta \psi = (\cdots) \phi \epsilon \end{cases}$ such that

- 1) leave invariant the SUSY action, i.e.  $\delta S_{green} = 0$
- 2) satisfy the "SUSY algebra", i.e.  $\begin{cases} [\delta_1, \delta_2] \phi = -2i(\bar{\epsilon}_2 \partial \epsilon_1) \phi \\ [\delta_1, \delta_2] \psi = -2i(\bar{\epsilon}_2 \partial \epsilon_1) \psi \end{cases}$

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 is called "SUSY parameter" so as if 
$$\left\{ \begin{array}{l} \partial_{\mu}\,\epsilon(x) = 0 & \text{global SUSY } \checkmark \\ \partial_{\mu}\,\epsilon(x) \neq 0 & \text{local SUSY} \end{array} \right.$$

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Naively, consider action principles of a real scalar and Majorana spinor fields

$$S = S_b + S_f = \frac{1}{2} \int d^4x \left[ (\partial_\mu \phi)(\partial^\mu \phi) + \bar{\psi}(i\partial\!\!\!/) \psi \right]$$

One can show that the above action is invariant under transformations

$$\begin{cases} \delta \phi &=& \bar{\epsilon} \, \psi \\ \delta \psi &=& -i \, \partial \!\!\!/ \phi \, \epsilon \end{cases}$$

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$$\delta_{1} \, \underline{\delta_{2}} \, \phi = \delta_{1} \, \overline{\epsilon_{2}} \, \psi$$

$$= \overline{\epsilon_{2}} \, \delta_{1} \psi$$

$$= \overline{\epsilon_{2}} \left( -i \, \partial \!\!\!/ \phi \, \epsilon_{1} \right) = -i \left( \overline{\epsilon_{2}} \, \partial \!\!\!/ \epsilon_{1} \right) \phi$$

$$[\delta_1, \delta_2] \phi = -i (\bar{\epsilon}_2 \partial \epsilon_1 - \bar{\epsilon}_1 \partial \epsilon_2) \phi = -2i (\bar{\epsilon}_2 \partial \epsilon_1) \phi \qquad \checkmark$$

$$[\delta_1, \delta_2] \psi = \cdots \qquad \neq -2i (\bar{\epsilon}_2 \not \partial \epsilon_1) \psi \qquad ??!!$$

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$$\begin{split} \delta_1 \, & \delta_2 \, \phi = \delta_1 \, \overline{\epsilon}_2 \, \psi \\ & = \overline{\epsilon}_2 \, \delta_1 \psi \\ & = \overline{\epsilon}_2 \left( -i \, \partial \!\!\!/ \phi \, \epsilon_1 \, \right) = -i \left( \, \overline{\epsilon}_2 \, \partial \!\!\!/ \epsilon_1 \, \right) \phi \end{split}$$

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# Example: chiral multiplet $(0, \tilde{0}; 1/2)$ :

$$S_{\scriptscriptstyle SUSY} = \, \tfrac{1}{2} \, \int d^4x \, \left[ \, (\partial_\mu \, \phi)(\partial^\mu \, \phi) \, \, + \, \, (\partial_\mu \, \tilde{\phi})(\partial^\mu \, \tilde{\phi}) \, \, + \, \, \bar{\psi} \, (\, i \! \partial \!\!\!/) \, \psi \, \, \right] \,$$

SUSY transformations 
$$\begin{cases} \delta \phi &= \bar{\epsilon} \, \psi \\ \delta \tilde{\phi} &= i \, \bar{\epsilon} \, \gamma^5 \, \psi \\ \delta \psi &= -i \, \partial \!\!\!/ \, \phi \, \epsilon - \gamma^5 \, \partial \!\!\!/ \, \bar{\phi} \, \epsilon \end{cases}$$

SUSY algebra 
$$\begin{cases} \left[\delta_{1}, \delta_{2}\right] \phi &= -2i\left(\bar{\epsilon}_{2} \not \!\!\! \partial \epsilon_{1}\right) \phi & \checkmark \\ \left[\delta_{1}, \delta_{2}\right] \tilde{\phi} &= -2i\left(\bar{\epsilon}_{2} \not \!\!\! \partial \epsilon_{1}\right) \tilde{\phi} & \checkmark \\ \left[\delta_{1}, \delta_{2}\right] \psi &\approx -2i\left(\bar{\epsilon}_{2} \not \!\!\! \partial \epsilon_{1}\right) \psi & \checkmark \end{cases}$$

where  $\approx$  denotes the equation of motion  $(i\partial \psi = 0)$  is used (on-shell supersymmetry)

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By defining the complex scalar field

$$\Phi = \frac{1}{\sqrt{2}} \left( \phi \, - \, i \, \tilde{\phi} \right) \qquad \Phi^\dagger = \frac{1}{\sqrt{2}} \left( \phi \, + \, i \, \tilde{\phi} \right)$$

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#### CSP

- is a massless elementary particle introduced by Wigner (1939)

$$\Phi(x,\eta) = \Phi(x) + \eta^{\mu} \Phi_{\mu}(x) + \frac{1}{2!} \eta^{\mu} \eta^{\nu} \Phi_{\mu\nu}(x) + \frac{1}{3!} \eta^{\mu} \eta^{\nu} \eta^{\rho} \Phi_{\mu\nu\rho}(x) \cdots$$

$$\Psi(x,\eta) = \Psi(x) + \eta^{\mu} \Psi_{\mu}(x) + \frac{1}{2!} \eta^{\mu} \eta^{\nu} \Psi_{\mu\nu}(x) + \frac{1}{3!} \eta^{\mu} \eta^{\nu} \eta^{\rho} \Psi_{\mu\nu\rho}(x) \cdots$$

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- 2 is a unitary irreducible representation (UIRs) of the Poincaré group
- $\odot$  characterizes by "continuous spin parameter"  $\mu$  (with the dimension of a mass
- 9 representation decomposes into a direct sum of all helicity reps., at  $\mu = 0$
- b has infinite degrees of freedom per space-time point
- wave equations were introduced by Wigner
- a has two types of representation; bosonic & fermionic

Bosonic CSP field

$$\Phi(x,\eta) = \Phi(x) + \eta^{\mu} \Phi_{\mu}(x) + \frac{1}{2!} \eta^{\mu} \eta^{\nu} \Phi_{\mu\nu}(x) + \frac{1}{3!} \eta^{\mu} \eta^{\nu} \eta^{\rho} \Phi_{\mu\nu\rho}(x) \cdots$$

Fermionic CSP field:

$$\Psi(x,\eta) = \Psi(x) + \eta^{\mu} \Psi_{\mu}(x) + \frac{1}{2!} \eta^{\mu} \eta^{\nu} \Psi_{\mu\nu}(x) + \frac{1}{3!} \eta^{\mu} \eta^{\nu} \eta^{\rho} \Psi_{\mu\nu\rho}(x) \cdots$$

§ states are NOT boost invariant (unlike the helicity states) while single-CSP is

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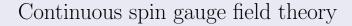
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Supersymmetric continuous spin gauge theory
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- In all SUSY theories, (# bosonic d.o.f) = (# fermionic d.o.f)
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- hence, in a CSP supermultiplet there would be four possibilities:

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 CSP supermultiplet  $\Rightarrow$  ( CSP , CSPino )  $\Downarrow$  ( real CSP , Majorana CSP)  $\times$  (complex CSP , Majorana CSP)  $\times$  ( real CSP , Dirac CSP )  $\times$  (complex CSP , Dirac CSP )  $\times$ 

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Unconstrained formalism of the SUSY CSP action is given by

$$S_{_{CSP}}^{^{SUSY}} = \left| S_{_{CSP}}^{\,b} + S_{_{CSP}}^{\,f} \right| = \int d^4x \, d^4\eta \left[ \Phi^\dagger(x,\eta) \;\; \mathbf{B} \;\; \Phi(x,\eta) \;\; + \;\; \overline{\Psi}(x,\eta) \;\; \mathbf{F} \;\; \Psi(x,\eta) \right]$$

where 
$$\begin{cases} \mathbf{B} = \delta'(\eta^2 + 1) \left[ -\Box + (\eta \cdot \partial_x)(\partial_\eta \cdot \partial_x + \mu) - \frac{1}{2}(\eta^2 + 1)(\partial_\eta \cdot \partial_x + \mu)^2 \right] \\ \mathbf{F} = \delta'(\eta^2 + 1) \left[ (\not \eta + i) \not \partial - (\eta^2 + 1)(\partial_\eta \cdot \partial_x + \mu) \right] \end{cases}$$

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$$\begin{cases} \delta \Phi(x,\eta) = \frac{1}{\sqrt{2}} \bar{\epsilon} \left( 1 + \gamma^5 \right) \left( \not \eta - i \right) \Psi(x,\eta) \\ \delta \Psi(x,\eta) = \frac{1}{\sqrt{2}} \left[ \not \partial - \frac{1}{2} \left( \not \eta + i \right) \left( \partial_{\eta} \cdot \partial_{x} + \mu \right) \right] \left( 1 - \gamma^5 \right) \epsilon \Phi(x,\eta) \end{cases}$$

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# Current work & Open problems

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• Supersymmetric unconstrained higher spin gauge theory in AdS<sub>4</sub> (arXiv: 20xx.xxxxx)

- group-theoretical meaning of CSP in AdS
- investigation of CSP in condensed matter systems
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Thank you for your attention!