



# Experimental Search for Higgs

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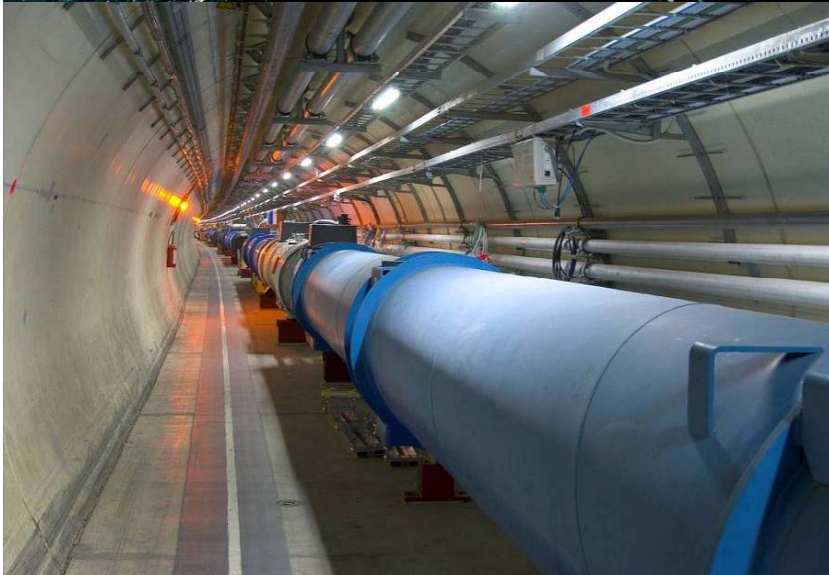
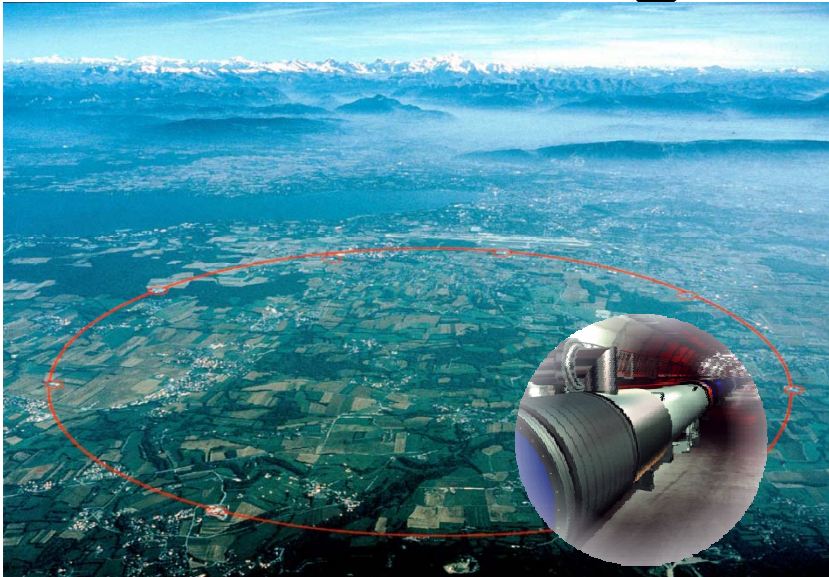
**One Day Workshop on the LHC Physics:  
Electroweak and Higgs**

**Aban 22, 1387**

# Outline

- LHC/CMS
- What we know now
- Higgs Production in LHC
- Higgs decay
- CMS/ATLAS potential for Higgs discovery

# The Large Hadron Collider



- 2835 + 2835 proton bunches separated by 7.5 m  
→ collisions every 25 ns  
= 40 MHz crossing rate
  - $10^{11}$  protons per bunch
  - at  $10^{34}/\text{cm}^2/\text{s}$   
≈ 25 pp interactions per crossing  
pile-up
  - ≈  $10^9$  pp interactions per second !!!
  - in each collision  
≈ 1600 charged particles produced
- enormous challenge for the detectors

# The LHC Machine and Experiments

## ATLAS/CMS Coverage

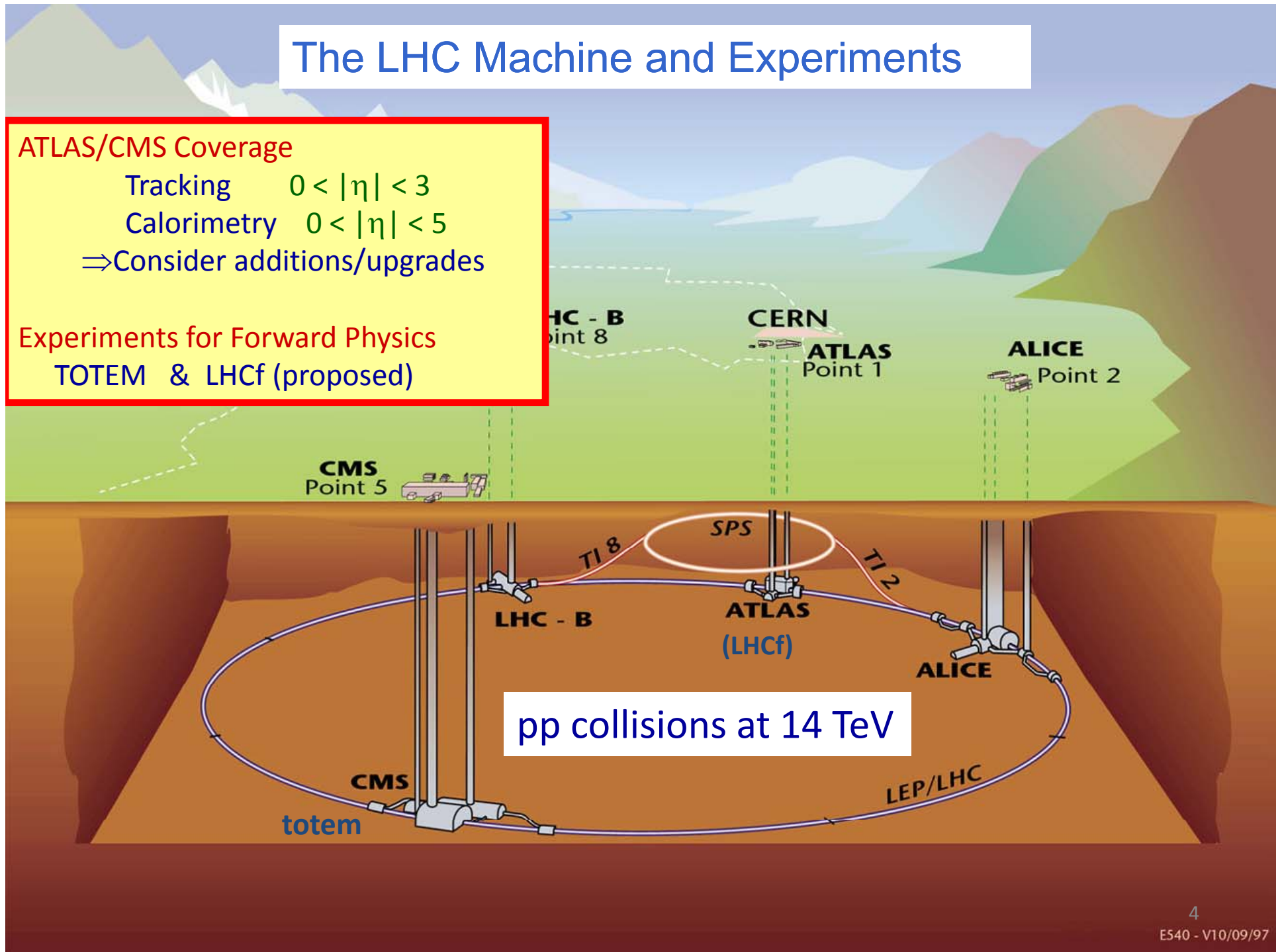
Tracking  $0 < |\eta| < 3$

Calorimetry  $0 < |\eta| < 5$

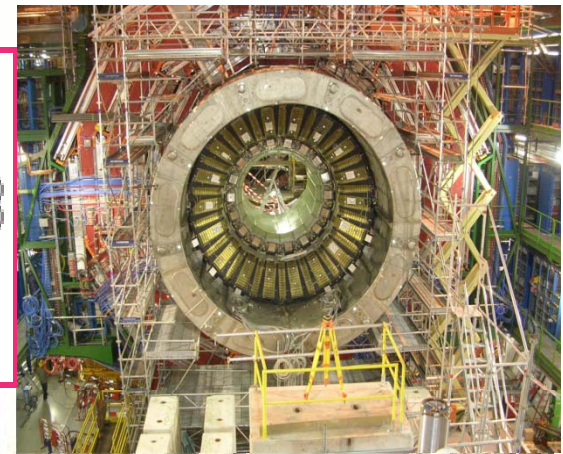
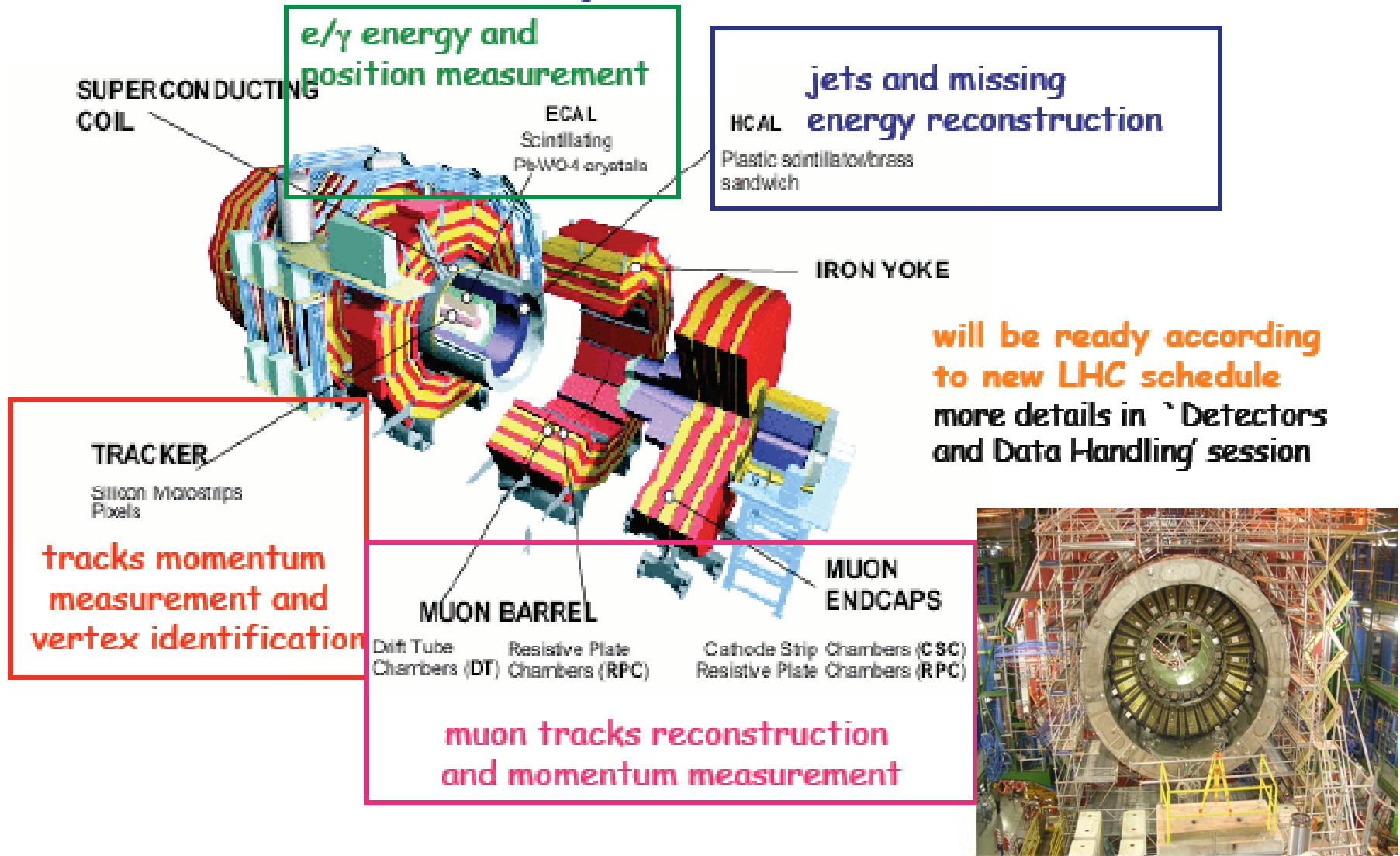
⇒ Consider additions/upgrades

## Experiments for Forward Physics

TOTEM & LHCf (proposed)



# CMS Experiment



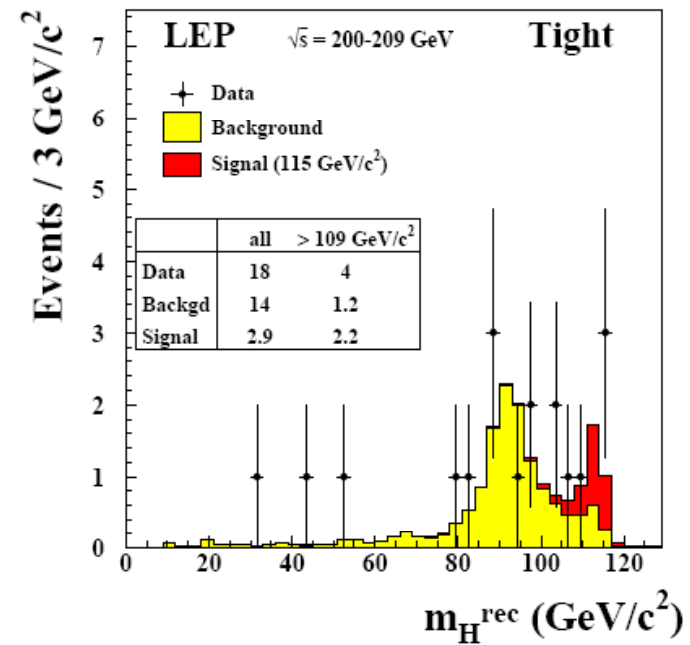
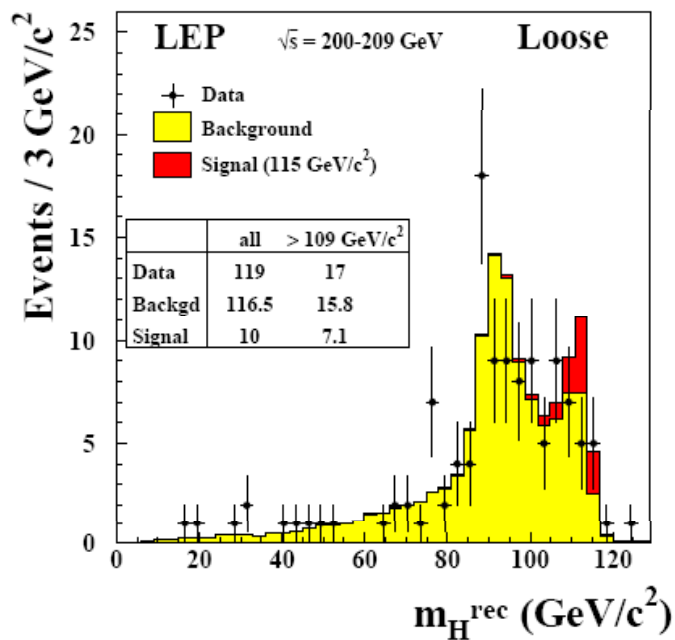
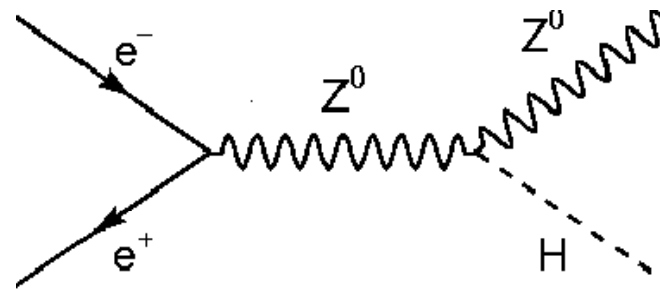
# CMS vs ATLAS

	ATLAS	CMS
MAGNET (S)	Air-core toroids + solenoid in inner cavity 4 magnets Calorimeters in field-free region	Solenoid Only 1 magnet Calorimeters inside field
TRACKER	Si pixels+ strips TRT → particle identification B=2T $\sigma/p_T \sim 5 \times 10^{-4} p_T \oplus 0.01$	Si pixels + strips No particle identification B=4T $\sigma/p_T \sim 1.5 \times 10^{-4} p_T \oplus 0.005$
EM CALO	Pb-liquid argon $\sigma/E \sim 10\%/ \sqrt{E}$ uniform longitudinal segmentation	PbWO <sub>4</sub> crystals $\sigma/E \sim 2-5\%/ \sqrt{E}$ no longitudinal segm.
HAD CALO	Fe-scint. + Cu-liquid argon (10 λ) $\sigma/E \sim 50\%/ \sqrt{E} \oplus 0.03$	Cu-scint. (> 5.8 λ +catcher) $\sigma/E \sim 100\%/ \sqrt{E} \oplus 0.05$
MUON	Air → $\sigma/p_T \sim 7\%$ at 1 TeV standalone	Fe → $\sigma/p_T \sim 5\%$ at 1 TeV combining with tracker

# Direct Search

Higgs-strahlung, the dominant production process for a standard model Higgs at LEP2.

direct (from LEP2):  $m_H > 114.4 \text{ GeV}/c^2$



## Constraints from precision data

$$\begin{aligned}\alpha &= \frac{1}{4\pi} \frac{g^2 g'^2}{g^2 + g'^2} = \frac{1}{137.03599976(50)} \\ G_F &= \frac{1}{\sqrt{2}v^2} = 1.16637(1) \times 10^{-5} \text{ GeV}^{-2} \\ m_Z &= \frac{1}{2} \sqrt{g^2 + g'^2} v = 91.1875(21) \text{ GeV} ,\end{aligned}$$

where the uncertainty is given in parentheses. The value of  $\alpha$  is extracted from **low-energy experiments**,  $G_F$  is extracted from the **muon lifetime**, and  $m_Z$  is measured from  **$e^+e^-$  annihilation** near the  $Z$  mass.

At tree level, we can express  $m_W$  as

$$m_W^2 = \frac{1}{\sin^2 \theta_W} \frac{\pi \alpha}{\sqrt{2} G_F}$$

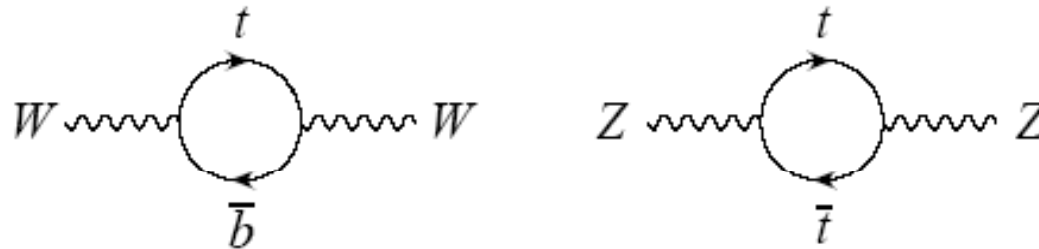
where

$$\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$$



## Clues to the Higgs boson mass

From the **sensitivity of electroweak observables** to the mass of the top, we are able to measure its mass, even **without directly producing it**



These quantum corrections alter the link between  $W$  and  $Z$  boson masses

$$m_W^2 = \frac{1}{\sin^2 \theta_W} \frac{\pi \alpha}{\sqrt{2} G_F} (1 - \Delta\rho)$$

$$\Delta\rho_{(\text{top})} \approx -\frac{3G_F}{8\pi^2\sqrt{2}} \frac{1}{\tan^2 \theta_W} m_t^2$$

The **strong dependence** on  $m_t^2$  accounts for the precision of the top-quark mass estimates derived from electroweak observables.

The Higgs boson quantum corrections are typically smaller than the top-quark corrections, and exhibit a more subtle dependence on  $m_H$  than the  $m_t^2$  dependence of the top-quark corrections.



$$\Delta\rho_{(\text{Higgs})} = \frac{11G_F m_Z^2 \cos^2 \theta_W}{24\sqrt{2}\pi^2} \log\left(\frac{m_H^2}{m_W^2}\right)$$

Since  $m_Z$  has been determined at LEP to 23 ppm, it is interesting to examine the dependence of  $m_W$  upon  $m_t$  and  $m_H$ .

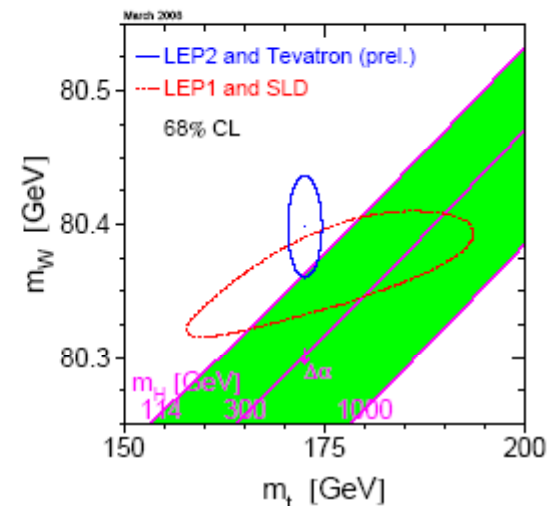
Indirect measurements of  $m_W$  and  $m_t$  (solid line)

Direct measurements of  $m_W$  and  $m_t$  (dotted line)

$$m_t = 172.6 \pm 1.4 \text{ GeV}$$

$$m_W = 80.398 \pm 0.025 \text{ GeV}$$

both shown as one-standard-deviation regions.



The indirect and direct determinations are in reasonable agreement and both favor a light Higgs boson, within the framework of the SM.

## SM Higgs mass fit to EW precision data

$$m_H = 87^{+36}_{-27} \text{ GeV}$$

Including theory uncertainty

$$m_H < 160 \text{ GeV} \quad (95\% \text{ CL})$$

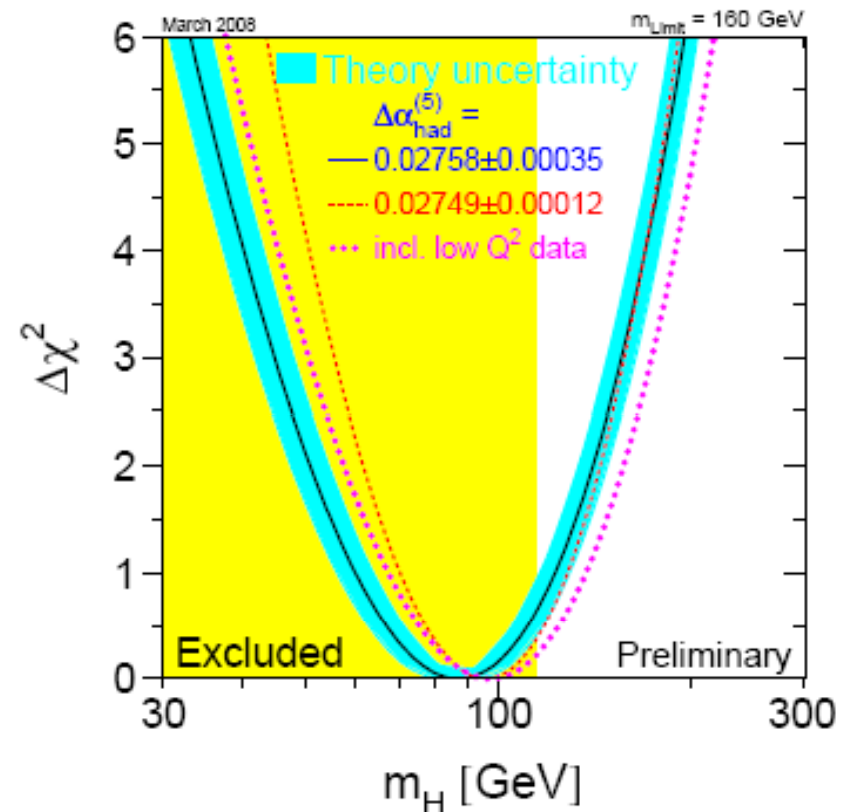
Does not include

Direct search limit from LEP

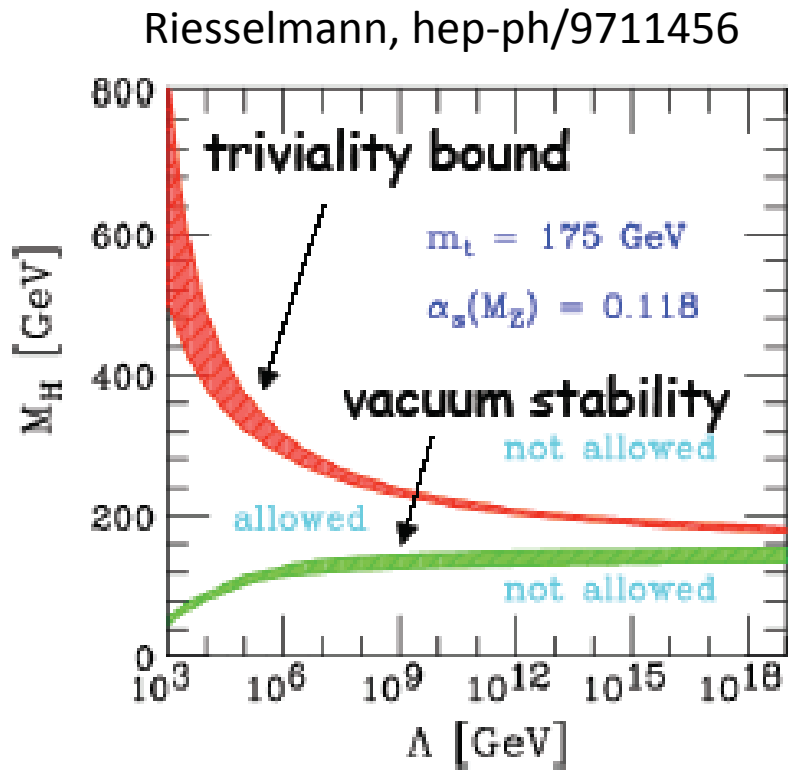
$$m_H > 114 \text{ GeV} \quad (95\% \text{ CL})$$

Renormalize probability for  
 $m_H > 114 \text{ GeV}$  to 100%:

$$m_H < 190 \text{ GeV} \quad (95\% \text{ CL})$$



# Theoretical Bounds

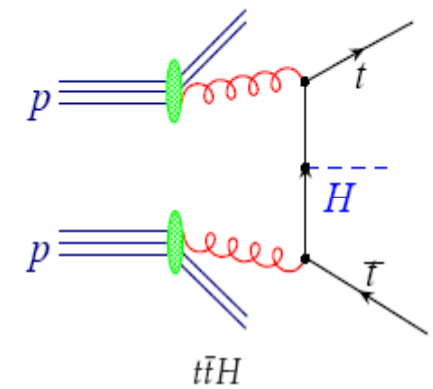
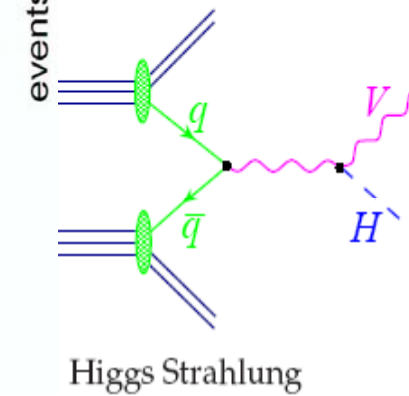
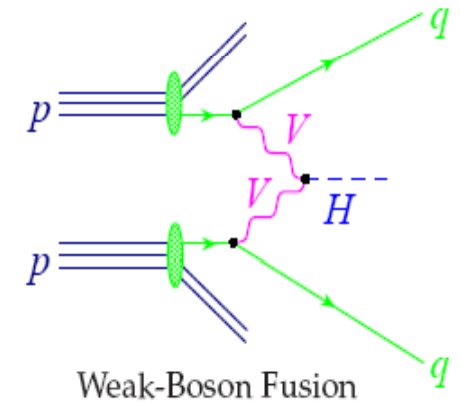
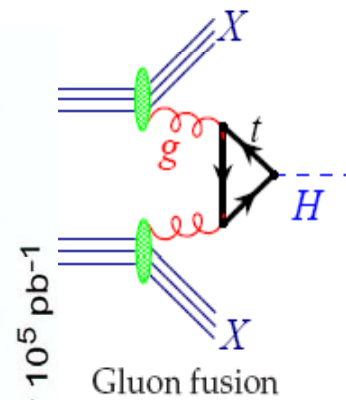
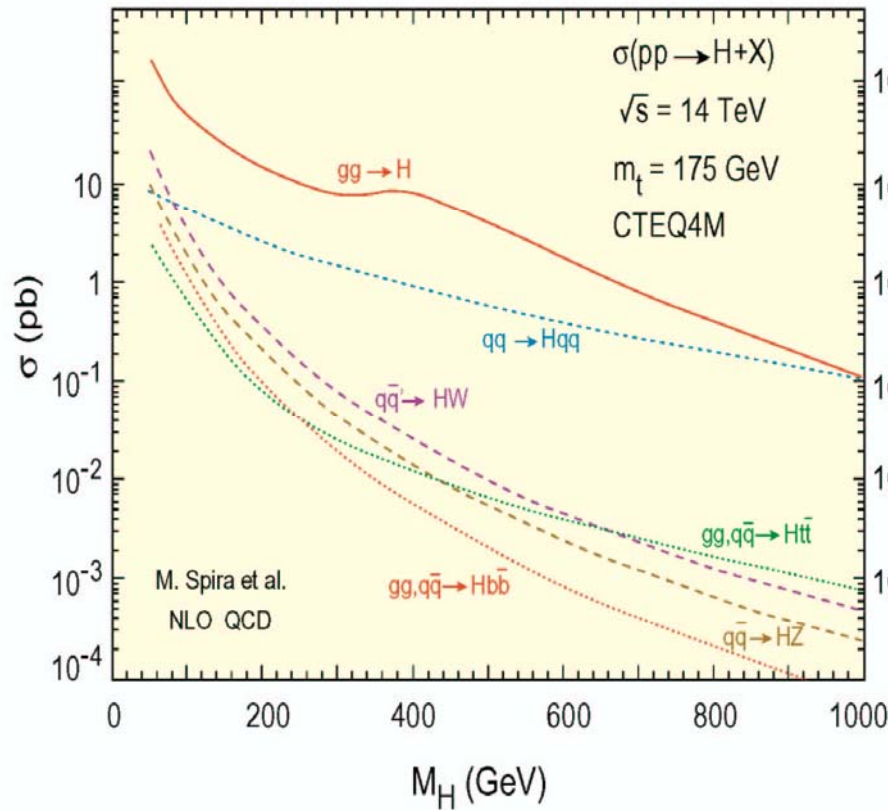


Notice the small window  
 $\sim 130 \text{ GeV} < M_H < \sim 180 \text{ GeV}$ , where  
 the theory is valid up to the Planck  
 scale

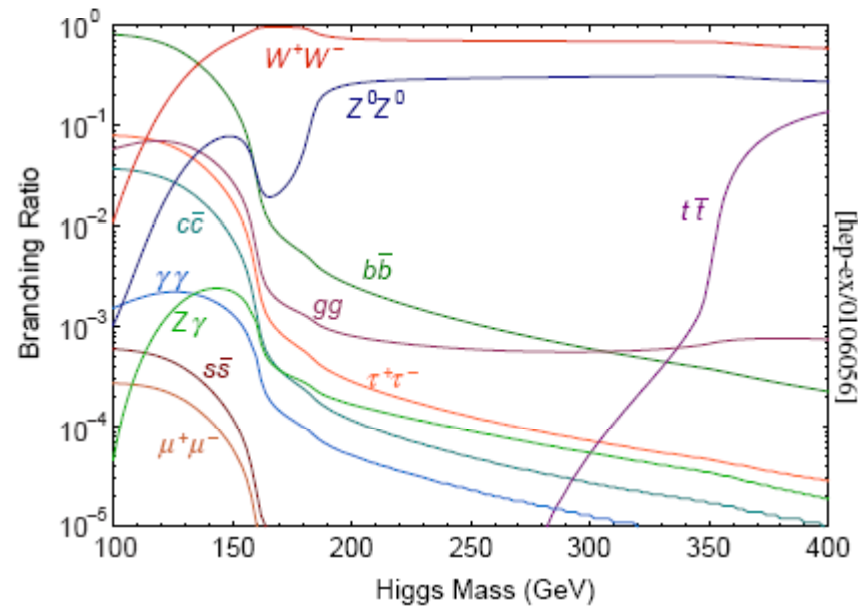
$$M_{Planck} = (\hbar c / G_{Newton})^{1/2} \approx 1.22 \times 10^{19} \text{ GeV}.$$

For a cutoff scale of  $\Lambda > 1000 \text{ TeV}$   
 the Higgs boson should lie in the mass  
 window  $110 \text{ GeV} < M_H < 300 \text{ GeV}$

# Higgs Production at LHC



# Higgs decay

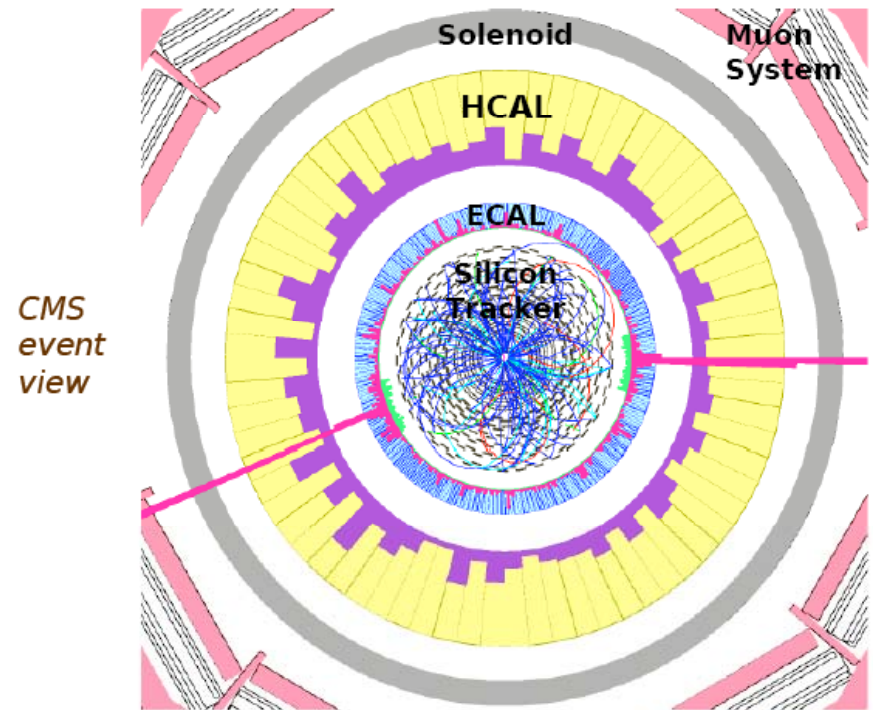
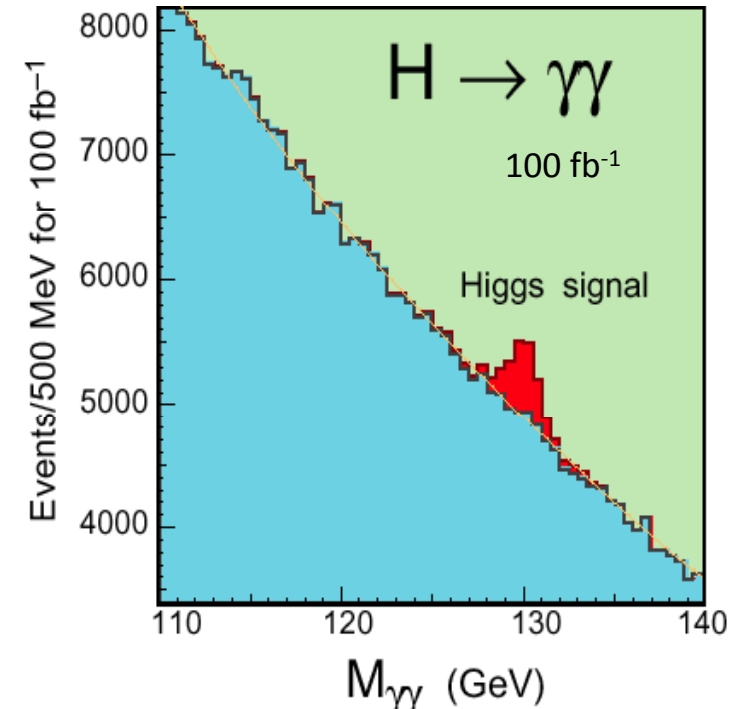


Above 350 GeV  $H \rightarrow t\bar{t}$  is opened,  
it is below  $ZZ$

# An intermediate mass Higgs ( $m_Z < m_H < 2m_Z$ )

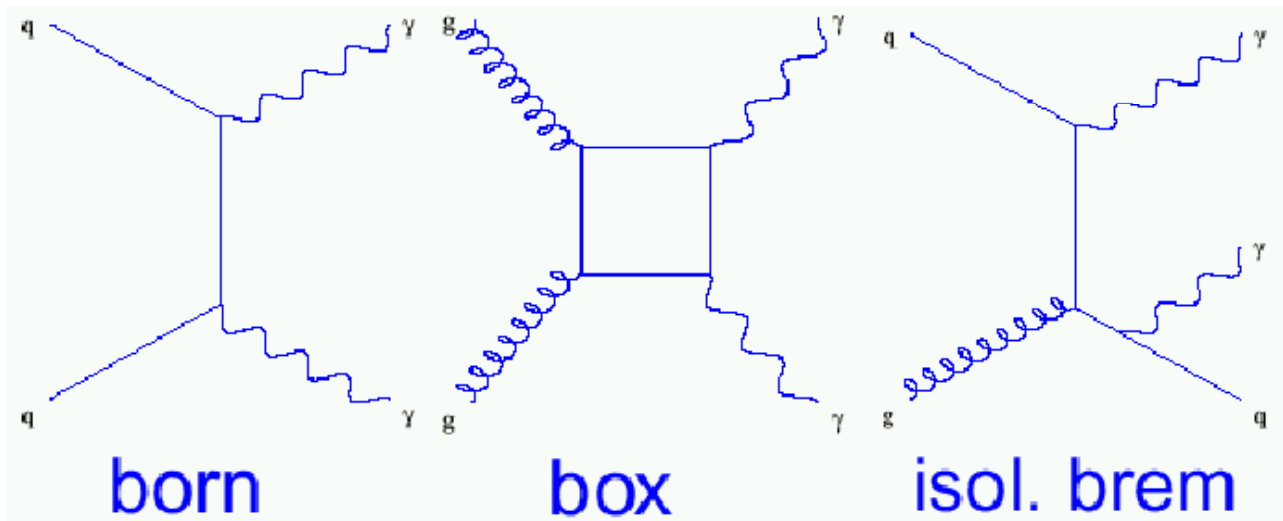
- Higgs below  $\sim 150$  GeV decays to  $b\text{-}\bar{b}$ , difficult to detect, QCD Bkg, low trigger efficiency (soft jets, no leptons).

- $H \rightarrow \gamma\gamma$  is much cleaner.  
**CMS is tuned for this range.**  
 BR  $\sim 0.002$  ( $M_H = 120$  GeV)



CMS NOTE - 2006/112

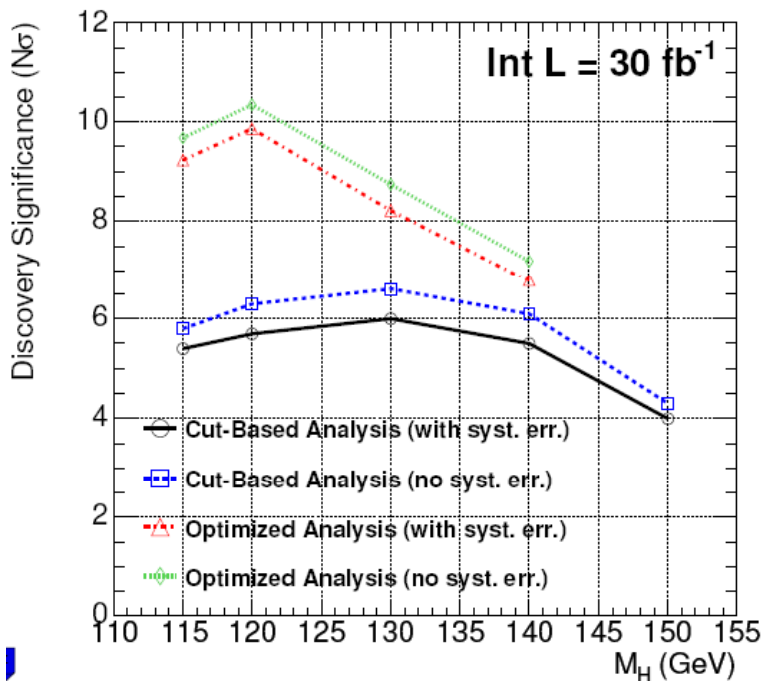
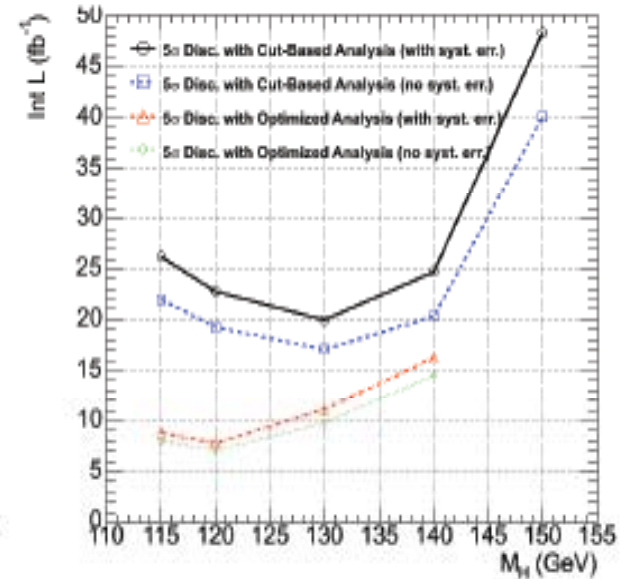
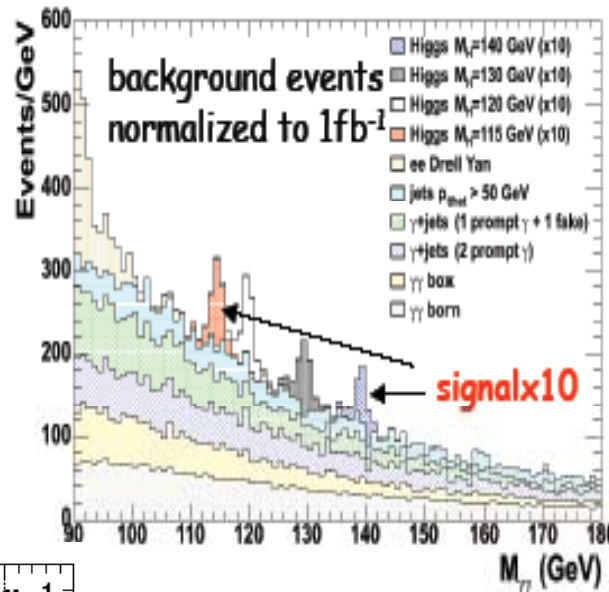
# Backgrounds for $H \rightarrow \gamma\gamma$



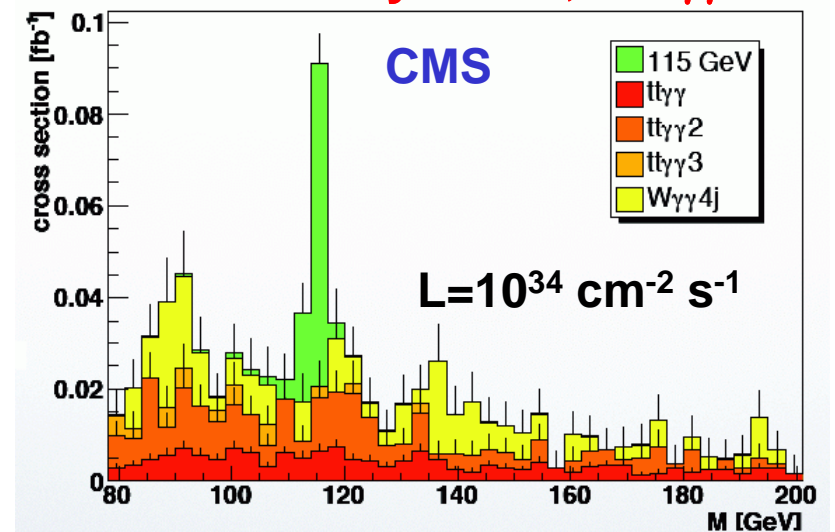
- important tool to suppress bkg.:  
photon ID and isolation: misidentified jets are accompanied by particles measured in tracker, e.m. and had. Calorimeters
- primary vertex:
  - e.m. energy measurement precise, but not the direction
  - longitudinal spread of interaction vertices  $\sim 53$  mm $\Rightarrow m_H$  smeared by 1.5 GeV if average vertex position is used



# Sensitivity for $H \rightarrow \gamma\gamma$



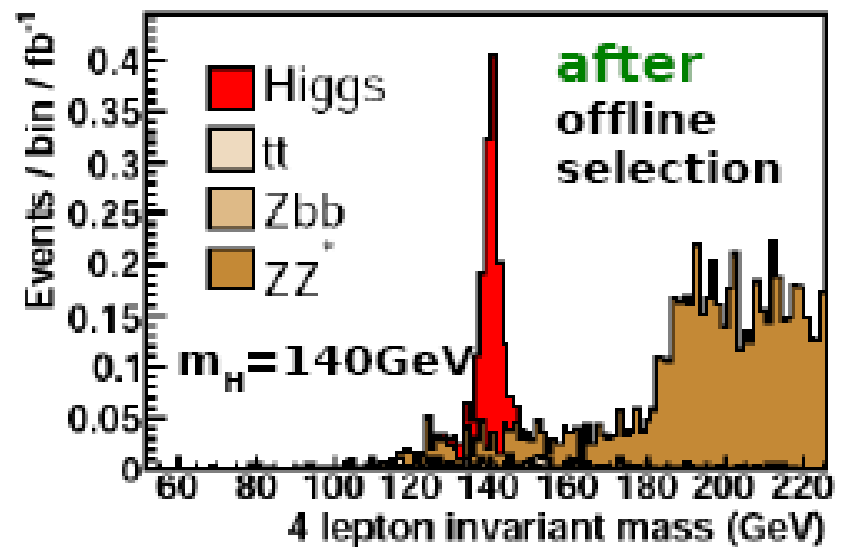
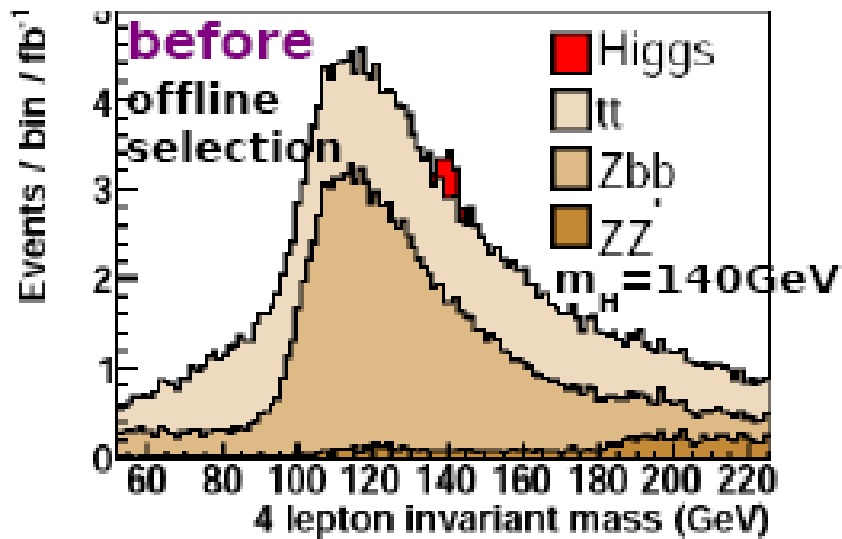
## Discovery of $t\bar{t}h$ , $h \rightarrow \gamma\gamma$



Higgs @ LHC  
Saeid Paktinat

# An intermediate mass higgs (cntd)

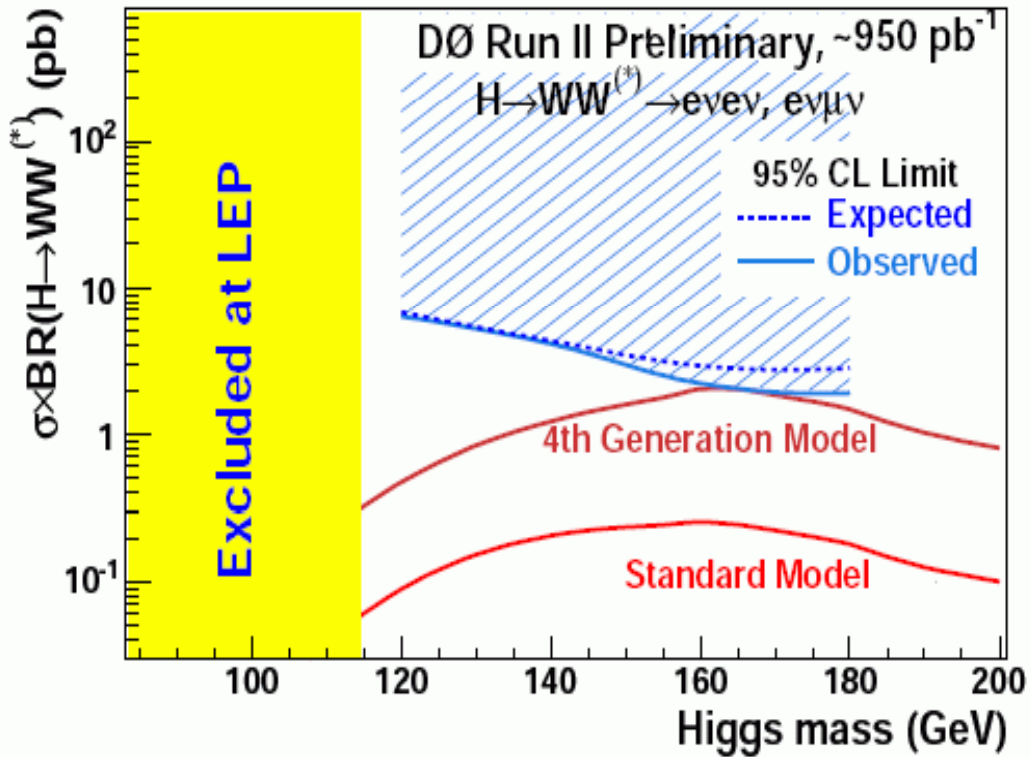
- $150 < M_H < 190$  GeV Thanks to clean leptons,  $H \rightarrow ZZ^*$  is the best channel to discover and characterize the Higgs. For  $Z^*$  the Z mass constraint does not work.



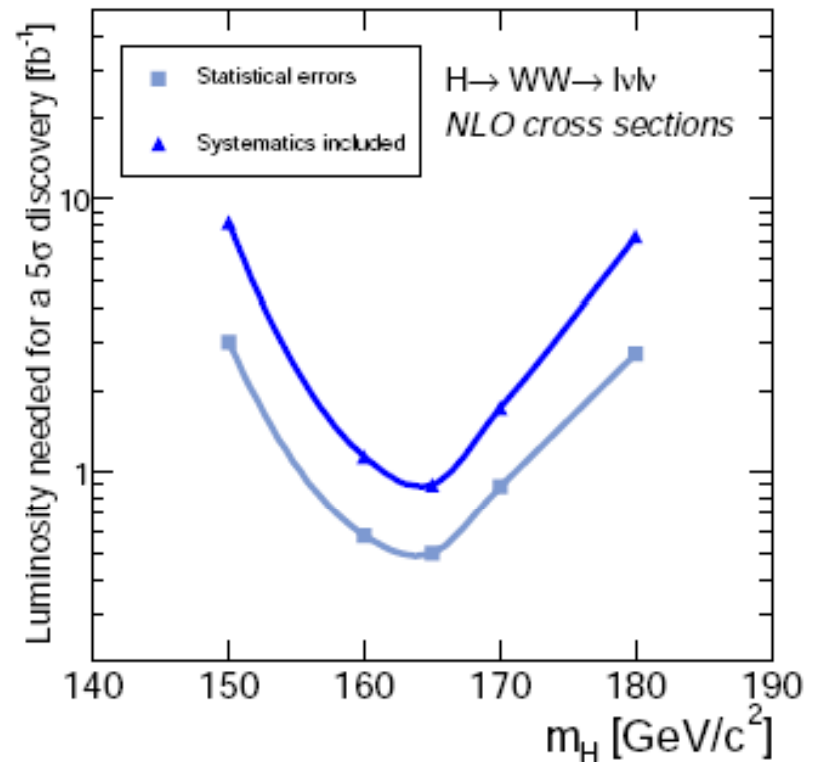
# Sensitivity for $H \rightarrow W^+W^- \rightarrow l\nu/l\nu$

## $150 < M_H < 190$ GeV

Direct exclusion from D0



CMS

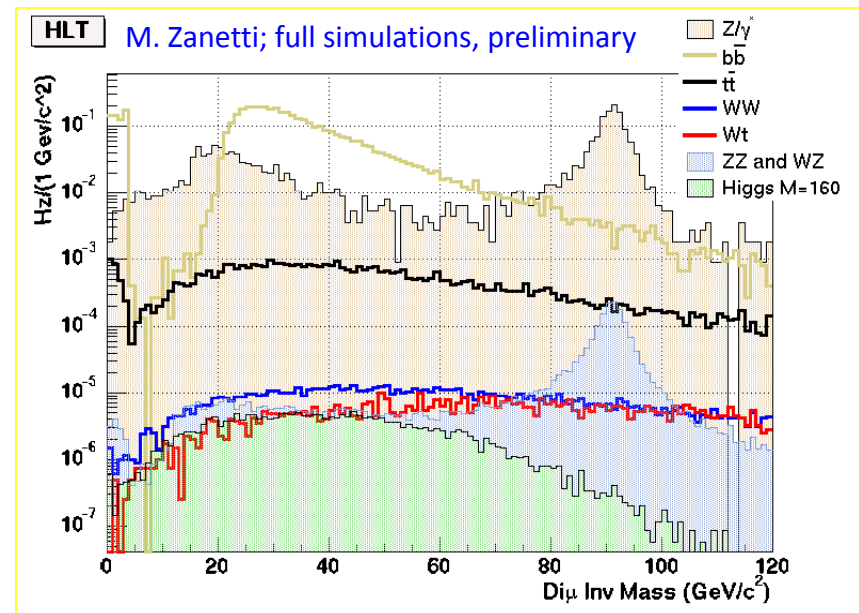
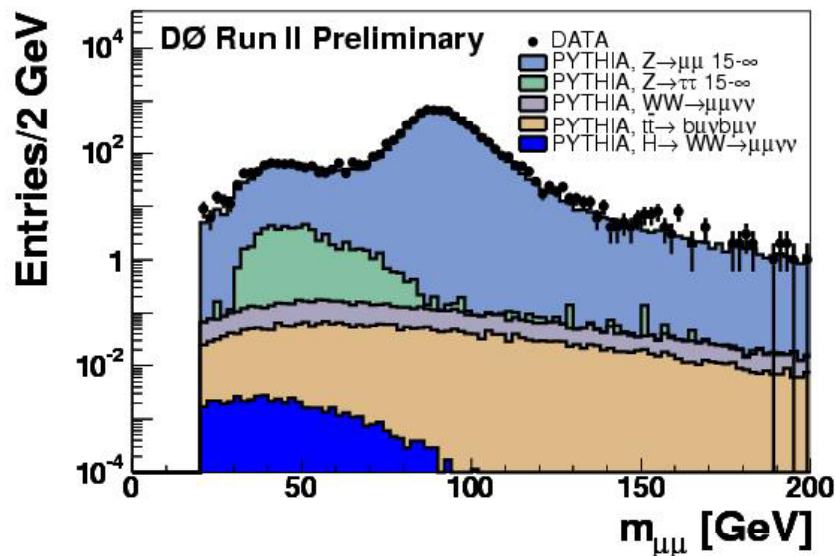


# H->WW->2l analysis at TeV and LHC (I)

Tevatron data and MC (PYTHIA)

LHC (CMS) Monte Carlo (PYTHIA)

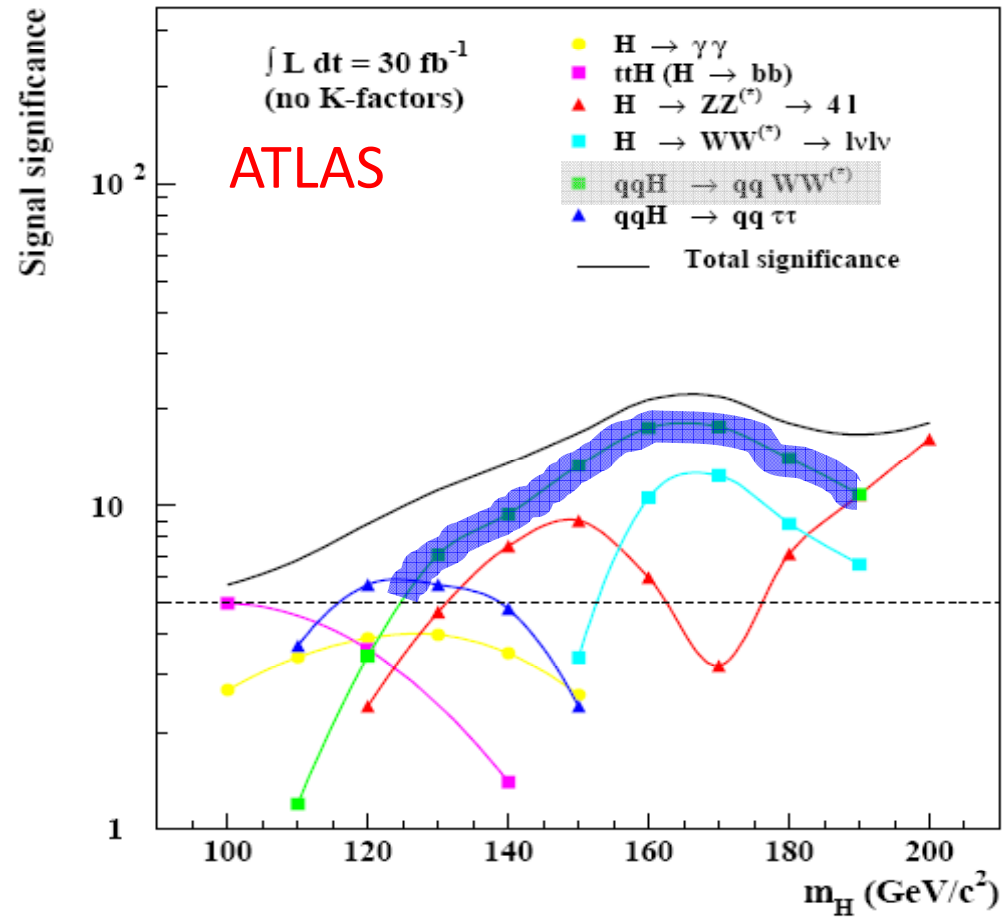
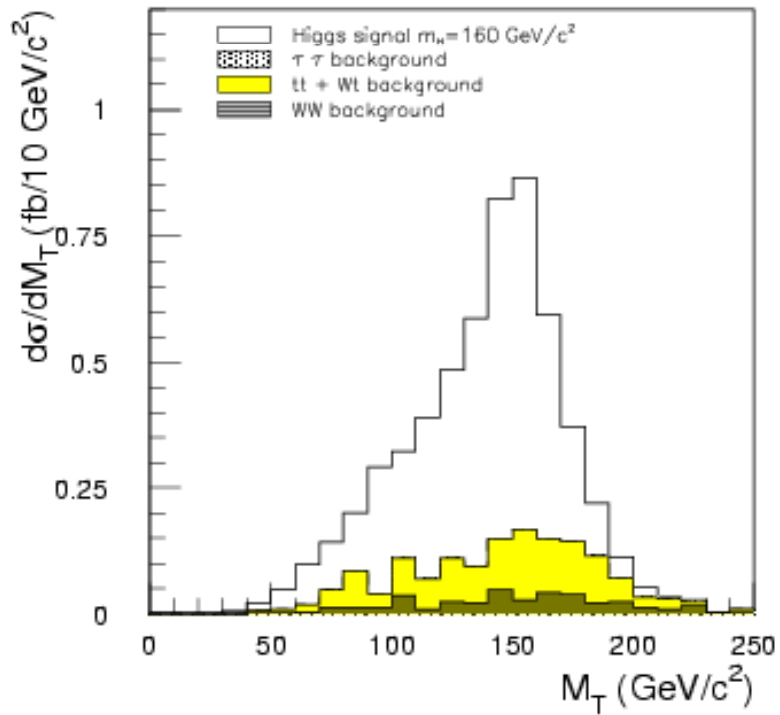
$M_H=160$  GeV



## Very similar event selections:

- cuts on lepton  $p_T$
- cut on miss  $E_T$ , Z resonance veto
- jet veto against  $t\bar{t}$
- $\Delta\phi(l\bar{l})$  cut is particularly important; exploit spin correlations

# ATLAS

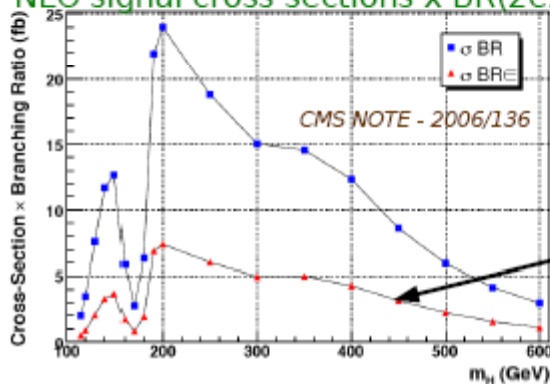


## A heavy Higgs ( $2m_Z < m_H < 650 \text{ GeV}$ )

- $H \rightarrow ZZ \rightarrow l^+l^-l^+l^-$  is the golden channel for Higgs decays.
- Irreducible Bkg is direct ZZ production. It can be controlled.
- With increasing the mass production/decay rate reduces e.g, less than 200 Higgs particles with  $m_H = 700 \text{ GeV}$  decay in  $H \rightarrow ZZ \rightarrow l^+l^-l^+l^-$  channel in a year at high luminosity.

# Sensitivity for $H \rightarrow ZZ \rightarrow 4l$

NLO signal cross sections x BR(2e2mu):

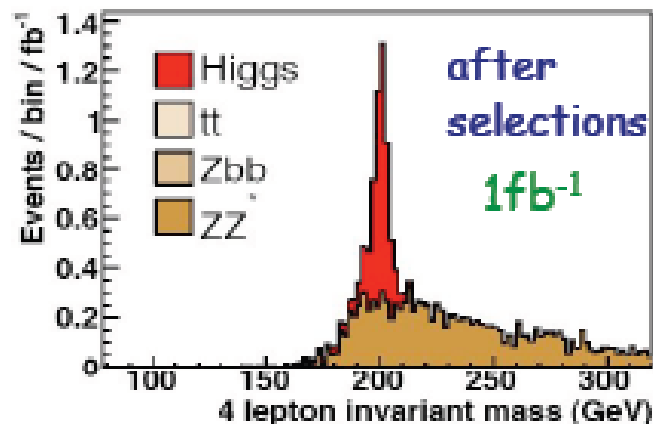
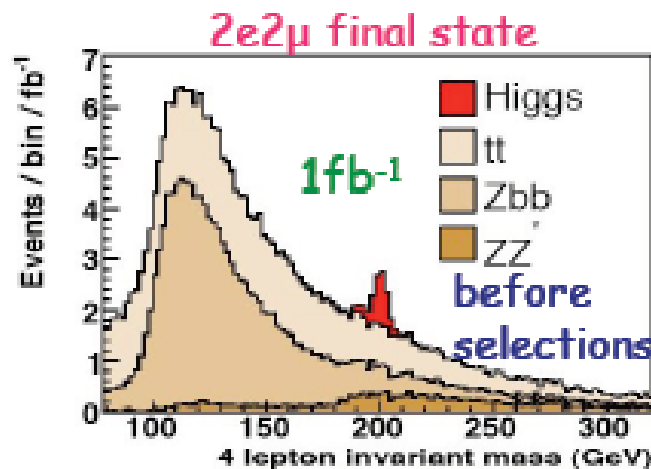
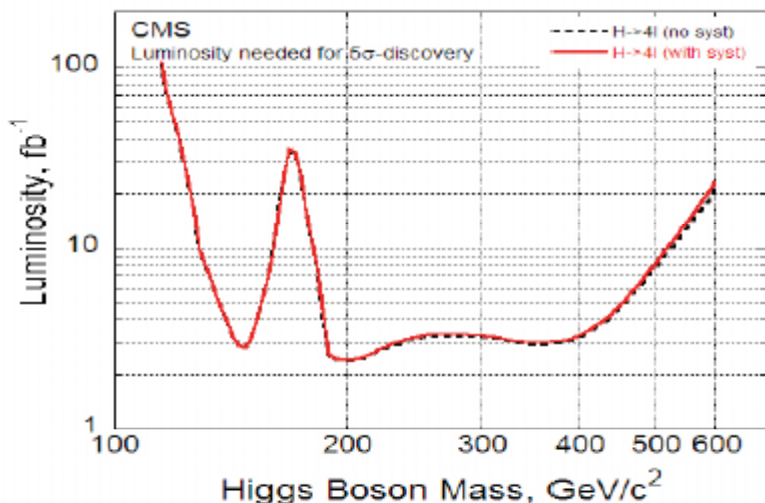


- gluon fusion and vector boson fusion considered

generator preselection: (CMS)

$p_t(e) > 5\text{GeV}$ ,  $\eta < 2.7$

$p_t(\mu) > 3\text{GeV}$ ,  $\eta < 2.4$



## A very heavy Higgs ( $m_H > 650$ GeV)

- To increase the statistics, the decay channel is changed.
- The decay channel  $H \rightarrow W^+W^- \rightarrow l\nu jj$ , has a branching ratio of just below 30% giving it a rate some 50 times higher than the four lepton channel from  $H \rightarrow ZZ$  decays.
- Mass measurement is very difficult and rough. No mass peak. Only event counting works.



# CMS capability

	H→bb	H→ττ	H→γγ	H→WW	H→ZZ
inclusive			YES	YES	YES
qqH		YES		YES	YES
W/Z+H					
ttH					

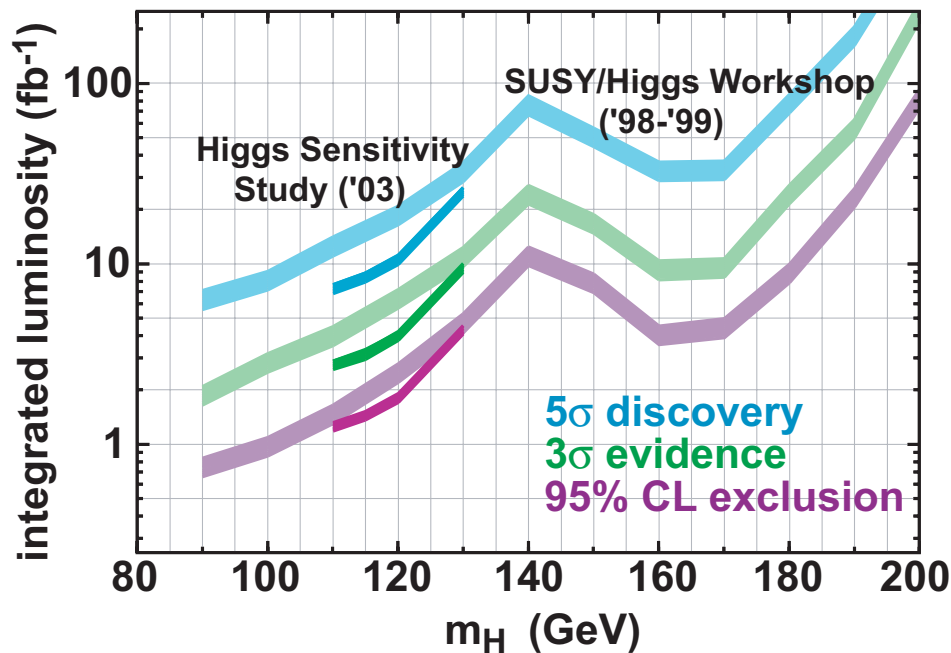
- ▷ filled boxes: detailed analysis available
- ▷ YES: sure discovery at  $\sqrt{L} < 30\text{fb}^{-1}$  in the appropriate masses range

# Prospects for SM Higgs

## Tevatron

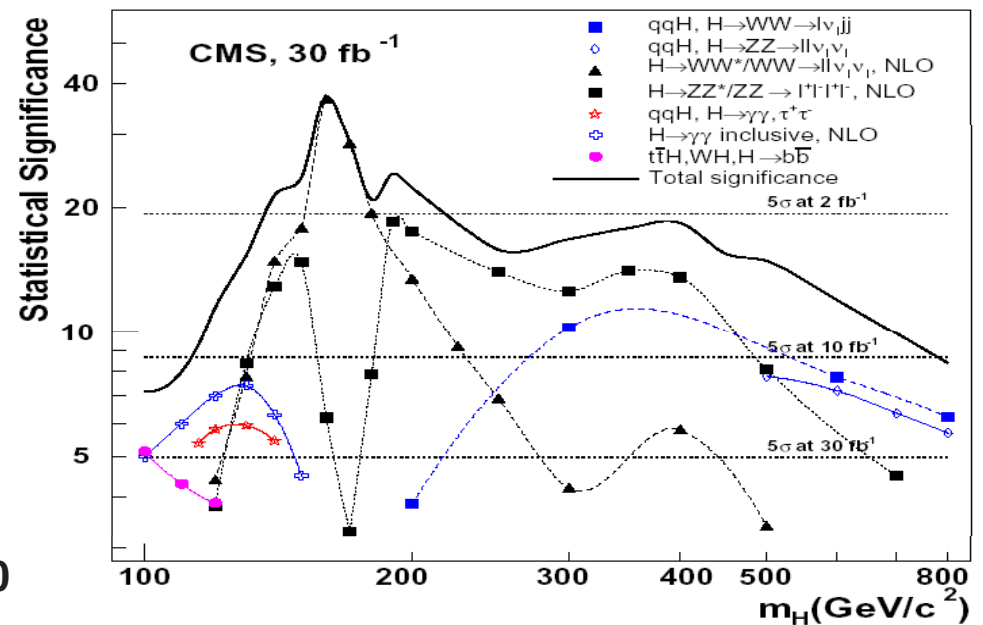
SM Higgs at 170 GeV/c<sup>2</sup>

- excluded at 95 % C.L. in 2008
- 3  $\sigma$  evidence for 120 GeV/c<sup>2</sup> in 2009 (?)

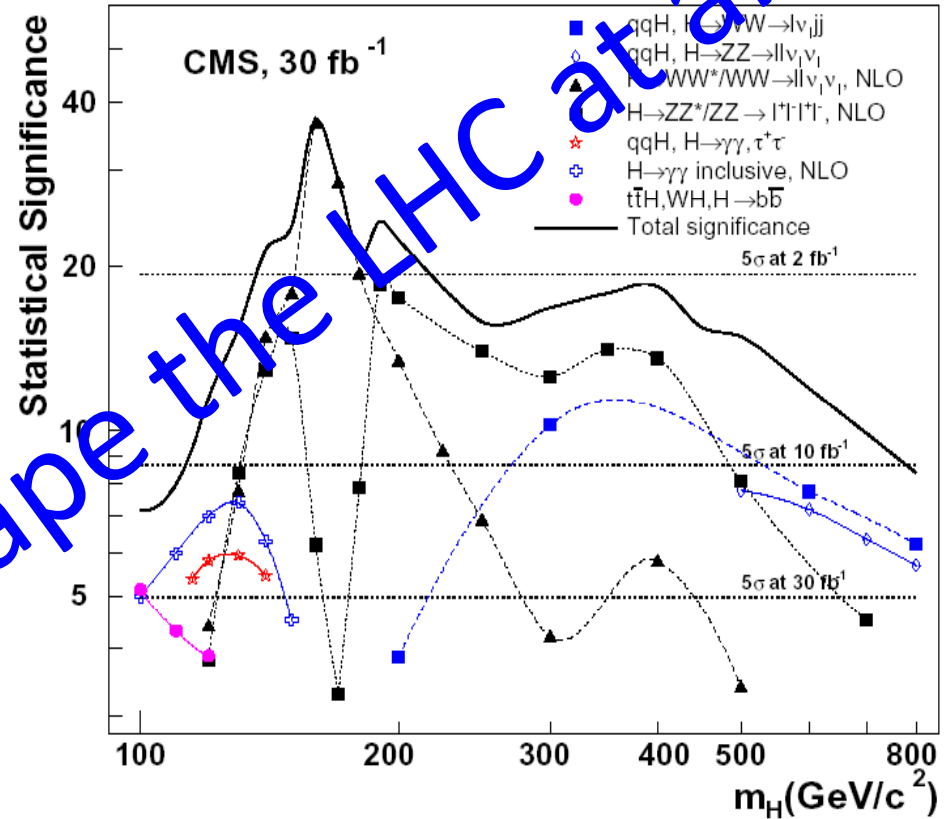
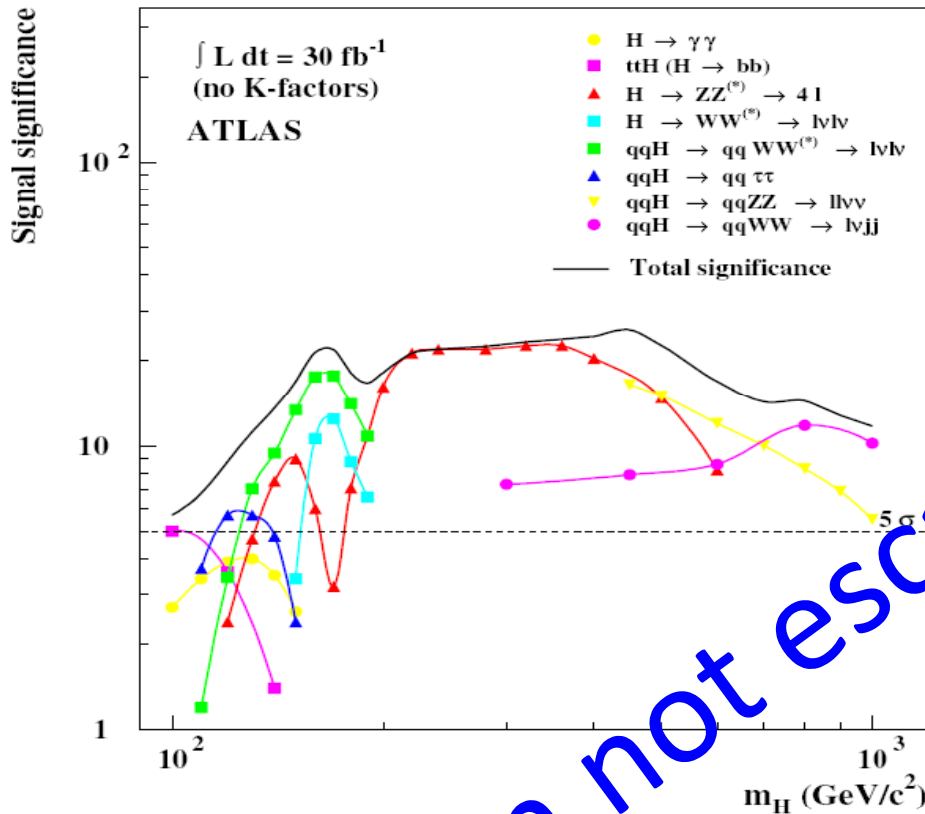


## LHC; one experiment (CMS example)

30 fb<sup>-1</sup> : 5  $\sigma$  discovery combining all channels for M<sub>H</sub> > 114 GeV;



# CMS vs ATLAS



Benchmark luminosities:

0.2 fb<sup>-1</sup>:

1 fb<sup>-1</sup>:

10 fb<sup>-1</sup>:

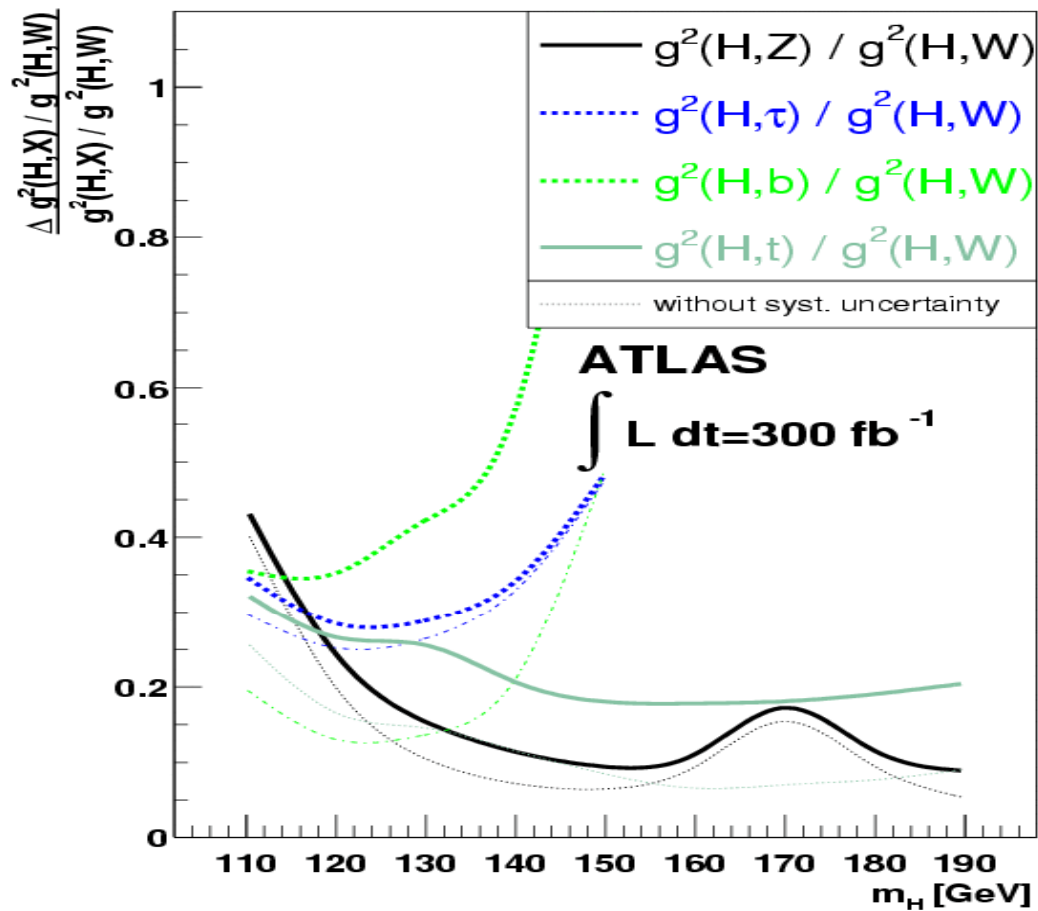
exclusion limits will start carving into SM Higgs x-section

discoveries become possible if  $M_H \sim 170 \text{ GeV}$

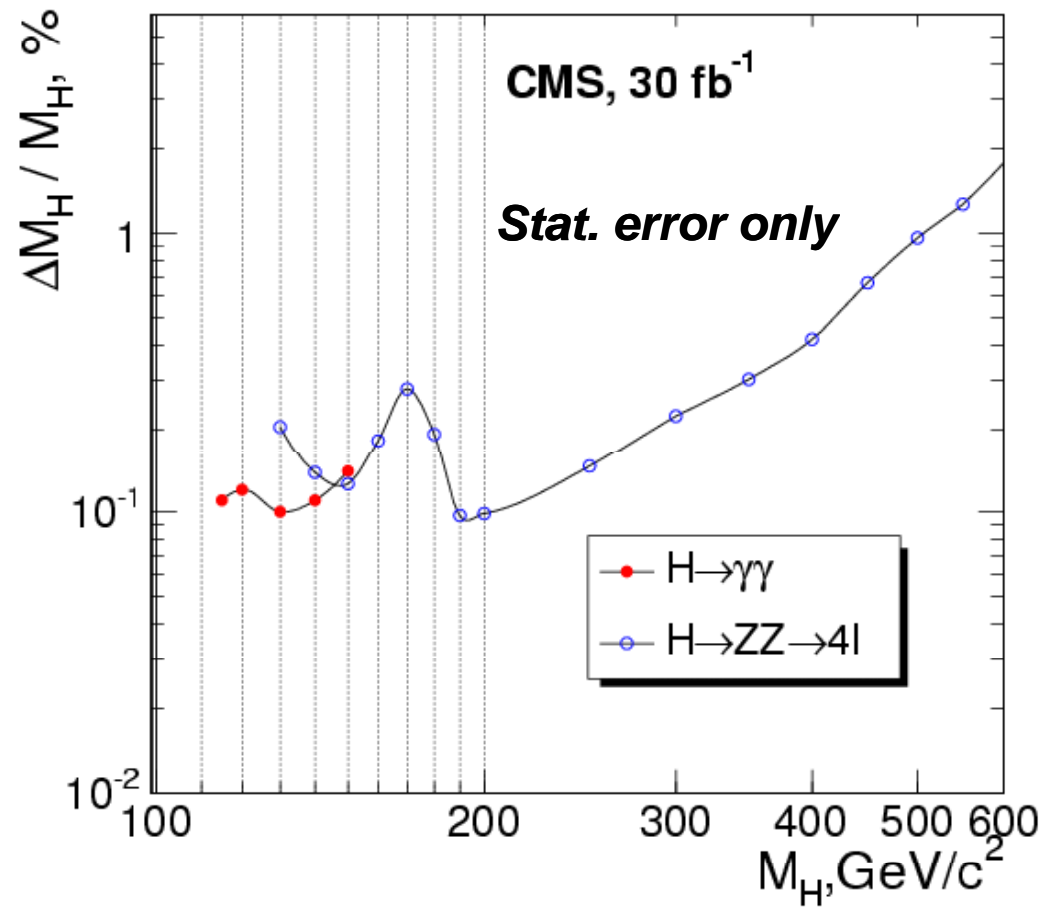
SM Higgs is discovered (or excluded) in full range

# Higgs Coupling measurement

ATL-PHYS-PUB-2003/030

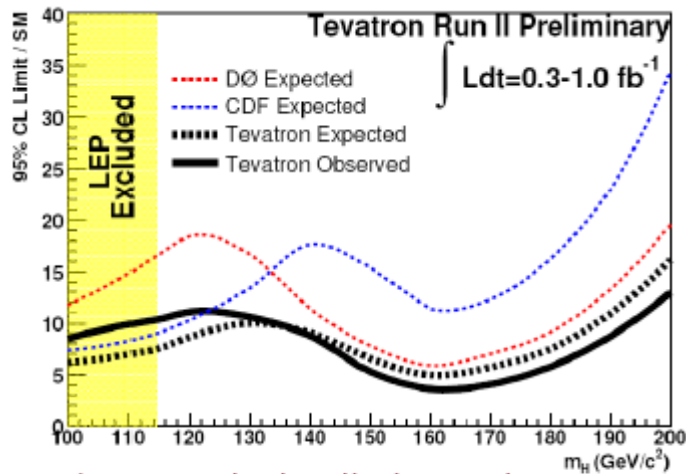


# Higgs Mass Accuracy



# Tevatron

search at Tevatron:  
most sensitive around  
160 GeV (H->WW)



(95% exclusion limit on Higgs  
production cross section  
normalized to SM cross section)

Summer 08 combined results of CDF and DØ  
H with 170 GeV/c<sup>2</sup> is excluded at 95 % C.L.

